Approximation of Schrödinger operators with δ -interactions supported on hypersurfaces

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We show that a Schrödinger operator $A_{\delta,\alpha}$ with a δ -interaction of strength α supported on a bounded or unbounded C^2 -hypersurface $\Sigma \subset \mathbb{R}^d$, $d \geq 2$, can be approximated in the norm resolvent sense by a family of Hamiltonians with suitably scaled regular potentials. The differential operator $A_{\delta,\alpha}$ with a singular interaction is regarded as a self-adjoint realization of the formal differential expression $-\Delta - \alpha \langle \delta_{\Sigma}, \cdot \rangle \delta_{\Sigma}$, where $\alpha \colon \Sigma \to \mathbb{R}$ is an arbitrary bounded measurable function. We discuss also some spectral consequences of this approximation result.

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1 Introduction

Singular Schrödinger operators with 'potentials' supported on subsets of the configuration space of a lower dimension are often used as models of physical systems because they are easier to solve, the original differential equation being reduced to the analysis of an algebraic or functional problem. The best known about them are solvable models with point interactions used in physics since the 1930s (see [26]), the rigorous analysis of which started from the seminal paper [8]; for a survey see the monograph [2]. In the last two decades the attention focused on interactions supported on curves, surfaces, and more complicated sets composed of them, which are used to model 'leaky' quantum systems in which the particle is confined to such manifolds or complexes, but the tunnelling between different parts of the interaction support is not neglected; for a review see [15] or [19, Chapter 10].

While these models are useful and mathematically accessible, one has to keep in mind that the singular interaction represents an idealized form of the actual, more realistic description. This naturally inspires the question about approximations of such singular potentials by regular ones. In the simplest case of a point interaction this problem was already addressed in the 1930s in [37]. Starting from the 1970s the approximation of Hamiltonians with point interactions supported on a finite or an infinite set of points in \mathbb{R}^d , $d \in \{1, 2, 3\}$, was treated systematically; cf. the monograph [2] and the references therein.

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Apart from that, the literature on the approximation of Schrödinger operators with δ -potentials supported on curves in \mathbb{R}^2 and surfaces in \mathbb{R}^3 is less complete; there are results available for the special cases that Σ is a sphere in \mathbb{R}^3 [5,34], that Σ is the boundary of a star-shaped domain in the plane [32], and that Σ is a smooth planar curve or surface and the interaction strength is constant [16,18]. In all of the above mentioned works convergence in the norm resolvent sense is shown. Abstract approaches developed in [3,35] cover more cases but imply only strong resolvent convergence in this context; for a very recent result on convergence in the operator norm in the case of compact hypersurfaces we refer to [22, Theorem 7.2]. We point out that the usage of scaled regular potentials is not the unique way of an approximation of δ -interactions supported on hypersurfaces, other mechanisms of approximation are discussed in e.g. [11, 12, 20, 31]. It is also worth mentioning that the approximation of δ -interactions supported on special periodic structures in \mathbb{R}^2 has important applications in the mathematical theory of photonic crystals, see [21] and the references therein.

The aim of the present paper is to analyze the general case where the interaction support is a C^2 -smooth hypersurface $\Sigma \subset \mathbb{R}^d$, $d \ge 2$, which is not necessarily bounded or closed, and the interaction strength is an arbitrary real valued bounded measurable function α on Σ . Following the approach of [16, 18] we show that the corresponding singular Schrödinger operator can be approximated in the norm resolvent sense by a family of regular ones with potentials suitably scaled in the direction perpendicular to Σ . We pay particular attention to the order of convergence and provide all preparatory technical integral estimates in a complete and self-contained form. We shall also mention some spectral consequences of the general approximation result.

In the following we describe our main result. Let $d \ge 2$ and let $\Sigma \subset \mathbb{R}^d$ be a bounded or unbounded orientable C^2 -hypersurface as in Definition 2.1, and consider the symmetric sesquilinear form

$$\mathfrak{a}_{\delta,\alpha}[f,g] = \left(\nabla f, \nabla g\right)_{L^2(\mathbb{R}^d;\mathbb{C}^d)} - \int_{\Sigma} \alpha \, f|_{\Sigma} \, \overline{g|_{\Sigma}} \, \mathrm{d}\sigma, \qquad \mathrm{dom} \, \mathfrak{a}_{\delta,\alpha} = H^1(\mathbb{R}^d),$$

where $\alpha \in L^{\infty}(\Sigma)$ is a real valued function and $f|_{\Sigma}$, $g|_{\Sigma}$ denote the traces of functions $f, g \in H^1(\mathbb{R}^d)$ on Σ . Standard arguments yield that $\mathfrak{a}_{\delta,\alpha}$ is a densely defined, closed, and semibounded form in $L^2(\mathbb{R}^d)$, and hence there exists a unique self-adjoint operator $A_{\delta,\alpha}$ in $L^2(\mathbb{R}^d)$ such that

$$(A_{\delta,\alpha}f,g) = \mathfrak{a}_{\delta,\alpha}[f,g], \qquad f \in \operatorname{dom} A_{\delta,\alpha}, \ g \in \operatorname{dom} \mathfrak{a}_{\delta,\alpha}, \tag{1.1}$$

see Lemma 2.7 for more details. The operator $A_{\delta,\alpha}$ is regarded as a Schrödinger operator with a δ -interaction of strength α supported on Σ which corresponds to the formal singular differential expression $-\Delta - \alpha \langle \delta_{\Sigma}, \cdot \rangle \delta_{\Sigma}$; cf. [10] and [7, Theorem 3.3]. The choice of the negative potential sign is motivated by the fact that interesting spectral features are in this context usually associated with attractive interactions.

Let ν be the continuous unit normal vector field on Σ , choose $\beta > 0$ sufficiently small as in Hypothesis 2.3 and consider layer neighborhoods Ω_{ε} of Σ of the form

$$\Omega_{\varepsilon} := \left\{ x_{\Sigma} + t\nu(x_{\Sigma}) : x_{\Sigma} \in \Sigma, \, t \in (-\varepsilon, \varepsilon) \right\}, \qquad 0 < \varepsilon \le \beta.$$

Fix a real valued potential $V \in L^{\infty}(\mathbb{R}^d)$ with support in Ω_{β} , define the scaled potentials $V_{\varepsilon} \in L^{\infty}(\mathbb{R}^d)$ with support in Ω_{ε} by

$$V_{\varepsilon}(x) := \begin{cases} \frac{\beta}{\varepsilon} V\left(x_{\Sigma} + \frac{\beta}{\varepsilon} t\nu(x_{\Sigma})\right), & \text{if } x = x_{\Sigma} + t\nu(x_{\Sigma}) \in \Omega_{\varepsilon}, \\ 0, & \text{else,} \end{cases}$$
(1.2)

and consider the corresponding self-adjoint Schrödinger operators

$$H_{\varepsilon}f = -\Delta f - V_{\varepsilon}f, \qquad \operatorname{dom} H_{\varepsilon} = H^{2}(\mathbb{R}^{d}).$$
(1.3)

With these preparatory considerations we can formulate the main result of the present paper.

Theorem 1.1 Let $\Sigma \subset \mathbb{R}^d$, $d \geq 2$, be an orientable C^2 -hypersurface as in Definition 2.1 which satisfies Hypothesis 2.3, let $Q \in L^{\infty}(\mathbb{R}^d)$ be real valued, and let $V \in L^{\infty}(\mathbb{R}^d)$ be real valued with support in Ω_{β} . Define $\alpha \in L^{\infty}(\Sigma)$ as the transversally averaged value of $t \mapsto V(x_{\Sigma} + t\nu(x_{\Sigma}))$ by

$$\alpha(x_{\Sigma}) := \int_{-\beta}^{\beta} V(x_{\Sigma} + s\nu(x_{\Sigma})) \mathrm{d}s \tag{1.4}$$

for a.e. $x_{\Sigma} \in \Sigma$ and let $A_{\delta,\alpha}$ be the corresponding Schrödinger operator with a δ -interaction of strength α supported on Σ . Then there exists a $\lambda_0 < 0$ such that $(-\infty, \lambda_0) \subset \rho(A_{\delta,\alpha} + Q) \cap \rho(H_{\varepsilon} + Q)$ for all $\varepsilon > 0$ sufficiently small and for every $\lambda \in (-\infty, \lambda_0)$ there is a constant $c = c(d, \lambda, \Sigma, V, Q) > 0$ such that

$$\left\| (H_{\varepsilon} + Q - \lambda)^{-1} - (A_{\delta,\alpha} + Q - \lambda)^{-1} \right\| \le c \varepsilon (1 + |\ln \varepsilon|)$$

holds for all $\varepsilon > 0$ sufficiently small. In particular, $H_{\varepsilon} + Q$ converges to $A_{\delta,\alpha} + Q$ in the norm resolvent sense, as $\varepsilon \to 0+$.

For given potentials Q and V Theorem 1.1 shows that the associated Hamiltonians $H_{\varepsilon} + Q$ in (1.3) converge in the norm resolvent sense to $A_{\delta,\alpha} + Q$, where the interaction strength α is as in (1.4). In the next corollary the converse situation is treated.

Corollary 1.2 Let $\Sigma \subset \mathbb{R}^d$, $d \geq 2$, be an orientable C^2 -hypersurface as in Definition 2.1 which satisfies Hypothesis 2.3, let $Q \in L^{\infty}(\mathbb{R}^d)$ be real valued, let $\alpha \in L^{\infty}(\Sigma)$ be real valued and let $A_{\delta,\alpha}$ be the corresponding Schrödinger operator with a δ -interaction of strength α supported on Σ . Define $V \in L^{\infty}(\mathbb{R}^d)$ by

$$V(x) := \begin{cases} \frac{1}{2\beta} \alpha(x_{\Sigma}), & \text{if } x = x_{\Sigma} + t\nu(x_{\Sigma}) \in \Omega_{\beta}, \\ 0, & \text{else,} \end{cases}$$

and let the associated scaled potentials V_{ε} and Hamiltonians H_{ε} be as in (1.2) and (1.3), respectively. Then there exists a $\lambda_0 < 0$ such that $(-\infty, \lambda_0) \subset \rho(A_{\delta, \alpha} + Q) \cap \rho(H_{\varepsilon} + Q)$ for all $\varepsilon > 0$ sufficiently small and for every $\lambda \in (-\infty, \lambda_0)$ there is a constant $c = c(d, \lambda, \Sigma, \alpha, Q) > 0$ such that

$$\left\| (H_{\varepsilon} + Q - \lambda)^{-1} - (A_{\delta,\alpha} + Q - \lambda)^{-1} \right\| \le c \varepsilon (1 + |\ln \varepsilon|)$$

holds for all $\varepsilon > 0$ sufficiently small. In particular, $H_{\varepsilon} + Q$ converges to $A_{\delta,\alpha} + Q$ in the norm resolvent sense, as $\varepsilon \to 0+$.

Let us briefly describe the structure of the proof of Theorem 1.1 and the contents of this paper. Section 2 contains preliminary material, definitions and properties of the hypersurfaces Σ and their layer neighborhoods, as well as a representation of the resolvent of the Schrödinger operator $A_{\delta,\alpha}$ which goes back to [10]. The heart of the proof of Theorem 1.1 is in Section 3. The main part of this section deals with the special case Q = 0. For this purpose, the potentials V_{ε} in (1.2) are factorized with the standard Birman-Schwinger method and useful representations of the resolvent of H_{ε} are provided. The convergence analysis is then done by comparing the different resolvent representations from Theorem 2.8 and Proposition 3.3, and essentially reduces to convergence properties of certain integral operators discussed in Lemma 3.4. However, the proof of Lemma 3.4 requires various refined technical estimates for integrals containing the Green's function for the free Laplacian which are outsourced to Appendix A. We wish to mention that in Appendix A particular attention is paid to keep the present paper self-contained. Therefore all necessary estimates are presented in full detail and complete rigorous form; as a result Appendix A is of mainly technical nature. The statement of Theorem 1.1 in the general case with $Q \neq 0$ follows then from the previous considerations by a simple perturbation argument. Eventually, there is a short Appendix B in which it is shown that boundaries of bounded C^2 -domains satisfy the assumptions imposed on the hypersurfaces Σ in this paper.

Finally, we agree that throughout the paper $c, C, C_k, \widetilde{C}_k, k \in \mathbb{N}$, denote constants that do not depend on space variables and on ε . In the formulation of the results we usually write C = C(...) to emphasize on which parameters these constants depend, but in the proofs we will mostly omit this.

2 Preliminaries

This section contains some preliminary material that will be useful in the main part of the paper. In Section 2.1 we recall certain basic facts from differential geometry of hypersurfaces in Euclidean spaces and of layers built around these hypersurfaces. Then, in Section 2.2 we define Schrödinger operators with δ -interactions supported on hypersurfaces in a mathematically rigorous way.

2.1 Hypersurfaces and their layer neighborhoods

In this section we introduce several notions associated to hypersurfaces and layers around these hypersurfaces. We follow the presentation from [27], which we adopt for our applications. We start with a suitable definition of a class of hypersurfaces in the Euclidean space \mathbb{R}^d . We wish to emphasize that the hypersurfaces considered here are in general unbounded and not necessarily closed; note also that the index set I in the parametrization below is assumed to be finite.

Definition 2.1 We call $\Sigma \subset \mathbb{R}^d$, $d \geq 2$, a C^2 -hypersurface and $\{\varphi_i, U_i, V_i\}_{i \in I}$ a parametrization of Σ , if I is a finite index set and the following holds:

- (a) $U_i \subset \mathbb{R}^{d-1}$ and $V_i \subset \mathbb{R}^d$ are open sets and $\varphi_i \colon U_i \to V_i$ is a C^2 -mapping for all $i \in I$;
- (b) rank $D\varphi_i(u) = d 1$ for all $u \in U_i$ and $i \in I$;
- (c) $\varphi_i(U_i) = V_i \cap \Sigma$ and $\varphi_i \colon U_i \to V_i \cap \Sigma$ is a homeomorphism;
- (d) $\Sigma \subset \bigcup_{i \in I} V_i;$
- (e) there exists a constant C > 0 such that

$$|\varphi_i(u) - \varphi_i(v)| \le C|u - v|$$

for all $u, v \in U_i$ and $i \in I$.

Let $\Sigma \subset \mathbb{R}^d$ be a C^2 -hypersurface with parametrization $\{\varphi_i, U_i, V_i\}_{i \in I}$. Then the inverse mappings φ_i^{-1} : $V_i \cap \Sigma \to U_i$ are often called charts and the family $\{\varphi_i^{-1}, V_i \cap \Sigma, U_i\}_{i \in I}$ atlas of Σ . For $x = \varphi_i(u) \in \Sigma$ $(u \in U_i, i \in I)$ we denote the tangent hyperplane by

$$T_x := \operatorname{span} \left\{ \partial_1 \varphi_i(u), \dots, \partial_{d-1} \varphi_i(u) \right\}.$$

The tangent hyperplane T_x is independent of the parametrization of Σ and dim $T_x = d - 1$ holds by Definition 2.1 (b). Subsequently, it is assumed that Σ is orientable, i.e. there exists a globally continuous unit normal vector field on Σ . From now on we fix a continuous unit normal vector field (which is unique up to multiplication with -1) and denote it by $\nu(x)$ for $x \in \Sigma$. Then the mapping $U_i \ni u \mapsto \nu(\varphi_i(u))$ is continuously differentiable for all $i \in I$ and $\partial_j \nu(\varphi_i(u)) \in T_{\varphi_i(u)}$ for all $u \in U_i$ and $j \in \{1, \ldots, d-1\}$, see, e.g., [27, Lemma 3.9 and Section 3F].

The first fundamental form I_x associated to Σ is the bilinear form on the tangent hyperplane T_x defined by

$$I_x[a,b] := \langle a,b \rangle, \qquad a,b \in T_x,$$

where $\langle \cdot, \cdot \rangle$ denotes the standard scalar product in \mathbb{R}^d . For $x = \varphi_i(u) \in \Sigma$ $(u \in U_i, i \in I)$ the matrix representing I_x in the canonical basis $\{(\partial_j \varphi_i)(u)\}_{j=1}^{d-1}$ of T_x is given by

$$G_i(u) = \left(\left\langle (\partial_k \varphi_i)(u), (\partial_l \varphi_i)(u) \right\rangle \right)_{k,l=1}^{d-1}$$
(2.1)

and also known as the metric tensor of Σ . Observe that $G_i(u) = (D\varphi_i(u))^\top \cdot (D\varphi_i(u))$. Together with condition (b) in Definition 2.1 this implies that $G_i(u)$ is positive definite.

Finally, we introduce the notion of the Weingarten map or shape operator.

Definition 2.2 Let $\Sigma \subset \mathbb{R}^d$, $d \geq 2$, be an orientable C^2 -hypersurface with parametrization $\{\varphi_i, U_i, V_i\}_{i \in I}$ and let $\nu(x), x \in \Sigma$, be a continuous unit normal vector field on Σ . For $x = \varphi_i(u) \in \Sigma$ ($u \in U_i, i \in I$) the Weingarten map $W(x) \colon T_x \to T_x$ is the linear operator acting on the basis vectors $\{\partial_j \varphi_i(u)\}_{j=1}^{d-1}$ of T_x as

$$W(x)\partial_j\varphi_i(u) := -\partial_j\nu(\varphi_i(u))$$

The Weingarten map W(x) is well defined (but its sign depends on the choice of the continuous unit normal vector field), independent of the parametrization and symmetric with respect to the inner product induced by the first fundamental form, see e.g. [27, Lemma 3.9] for the case d = 3. For $x = \varphi_i(u) \in \Sigma$ ($u \in U_i$, $i \in I$) the matrix associated to the linear mapping W(x) corresponding to the canonical basis $\{\partial_j \varphi_i(u)\}_{j=1}^{d-1}$ of T_x will be denoted by $L_i(u)$.

The eigenvalues $\{\mu_j(u)\}_{j=1}^{d-1}$ of $L_i(u)$ are the principal curvatures of Σ and do not depend on the choice of the parametrization, see [27, Definition 3.46]. In particular, the quantity $\det(1-tL_i(u))$ for $t \in \mathbb{R}$, which will appear later frequently, is independent of the parametrization and will also be denoted by $\det(1-tW(x))$. Furthermore, the eigenvalues of W(x) depend continuously on $x \in \Sigma$, as the entries of L_i depend continuously on $u \in U_i$ (see the text after Definition 3.10 in [27]) and $\varphi_i : U_i \to \Sigma \cap V_i$ is a homeomorphism.

Next, we discuss a convenient definition of an integral for functions defined on the C^2 -hypersurface Σ . For this fix a parametrization $\{\varphi_i, U_i, V_i\}_{i \in I}$ of Σ as in Definition 2.1 with a finite index set I and choose a partition of unity subordinate to $\{V_i\}_{i \in I}$, that is a family of functions $\chi_i : \mathbb{R}^d \to [0, 1], i \in I$, with the following properties:

- (i) $\chi_i \in C^{\infty}(\mathbb{R}^d)$ for all $i \in I$;
- (ii) supp $\chi_i \subset V_i$ for all $i \in I$;
- (iii) $\sum_{i \in I} \chi_i(x) = 1$ for any $x \in \Sigma$.

Note that some of the functions χ_i , $i \in I$, are not compactly supported, if Σ is unbounded.

A function $f: \Sigma \to \mathbb{C}$ is said to be measurable (integrable), if

$$U_i \ni u \mapsto \chi_i(\varphi_i(u)) f(\varphi_i(u))$$

is measurable (integrable, respectively) for all $i \in I$. If $f: \Sigma \to \mathbb{C}$ is integrable, we define the integral of f over Σ as

$$\int_{\Sigma} f(x) \mathrm{d}\sigma(x) := \sum_{i \in I} \int_{U_i} \chi_i(\varphi_i(u)) f(\varphi_i(u)) \sqrt{\det G_i(u)} \mathrm{d}u,$$
(2.2)

where $du := d\Lambda_{d-1}(u)$ denotes the usual (d-1)-dimensional Lebesgue measure on U_i and $G_i(u)$ is the matrix of the first fundamental form given in (2.1). The measure σ in (2.2) coincides with the canonical Hausdorff measure on Σ which is independent of the parametrization of Σ ; cf. [28, Appendix C.8]. Therefore, the above definition of the integral does not depend on the parametrization of Σ and the choice of the partition of unity. We denote the space of (equivalence classes of) square integrable functions $f : \Sigma \to \mathbb{C}$ with respect to σ by $L^2(\Sigma)$.

Next, we introduce layer neighborhoods of a C^2 -hypersurface Σ and we impose some additional conditions on Σ in Hypothesis 2.3 below. For this it is useful to define the functions

$$\iota_{\varphi_i} \colon U_i \times \mathbb{R} \to \mathbb{R}^d, \qquad \iota_{\varphi_i}(u, t) := \varphi_i(u) + t\nu(\varphi_i(u)), \qquad i \in I.$$
(2.3)

The Jacobian matrix of ι_{φ_i} , $i \in I$, is given by the $d \times d$ matrix

$$(D\iota_{\varphi_i})(u,t) = \begin{pmatrix} (D\varphi_i)(u)(1-tL_i(u)) & \nu(\varphi_i(u)) \end{pmatrix}$$

and the absolute value of the determinant of this matrix can be expressed as

$$\left|\det\left((D\iota_{\varphi_i})(u,t)\right)\right| = \left|\det(1 - tL_i(u))\right| \sqrt{\det G_i(u)};$$
(2.4)

cf. [29, Section 2] and [14, Section 3]. We will also make use of the mapping

$$\iota_{\Sigma} \colon \Sigma \times \mathbb{R} \to \mathbb{R}^{d}, \qquad \iota_{\Sigma}(x_{\Sigma}, t) := x_{\Sigma} + t\nu(x_{\Sigma}), \tag{2.5}$$

and layer neighborhoods Ω_β of Σ of the form

$$\Omega_{\beta} := \iota_{\Sigma}(\Sigma \times (-\beta, \beta)), \qquad \beta > 0.$$
(2.6)

We employ the following hypothesis for the hypersurface Σ .

Hypothesis 2.3 Let $\Sigma \subset \mathbb{R}^d$ be an orientable C^2 -hypersurface with parametrization $\{\varphi_i, U_i, V_i\}_{i \in I}$. Assume that there exists $\beta > 0$ such that

- (a) the restriction of the mapping ι_{Σ} on $\Sigma \times (-\beta, \beta)$ is injective;
- (b) there is a constant $\eta \in (0,1)$ such that $\det(1-tW(x_{\Sigma})) \in (1-\eta, 1+\eta)$ for all $x_{\Sigma} \in \Sigma$ and all $t \in (-\beta, \beta)$;
- (c) there exists a constant c > 0 such that the mappings ι_{φ_i} in (2.3) satisfy

$$|\iota_{\varphi_i}(u,t) - \iota_{\varphi_i}(v,s)|^2 \ge c^2 (|u-v|^2 + |s-t|^2)$$

for all $u, v \in U_i$, $s, t \in (-\beta, \beta)$ and all $i \in I$.

All the assumptions of Hypothesis 2.3 are satisfied for the boundary of a compact and simply connected C^2 domain; see Appendix B. We also mention that a similar set of assumptions was imposed in [10, Section 4]. In the next proposition it will be shown that item (c) in Hypothesis 2.3 implies that the eigenvalues of W are uniformly bounded on Σ . In particular, this shows that (b) in Hypothesis 2.3 is then automatically satisfied if $\beta > 0$ is small enough.

Proposition 2.4 Let $\Sigma \subset \mathbb{R}^d$ be an orientable C^2 -hypersurface and assume that item (c) in Hypothesis 2.3 holds. Then the eigenvalues of the matrix of the Weingarten map are uniformly bounded on Σ .

Proof. Let $\beta > 0$ be as in Hypothesis 2.3 (c) and suppose that the eigenvalues of the Weingarten map are not uniformly bounded. Then for some $i \in I$ there exists $u \in U_i$ and an eigenvalue μ of $L_i(u)$ such that $|\mu| > \beta^{-1}$. Choose a sequence $(s_n) \subset (-\beta, \beta)$ such that s_n^{-1} are not eigenvalues of $L_i(u)$ and $s_n \to \mu^{-1}$. Then

$$\det(1 - s_n L_i(u)) \neq 0$$
 and $\det(1 - s_n L_i(u)) \rightarrow 0$

and as $G_i(u)$ is positive definite and has uniformly bounded values by Definition 2.1 (e), the same holds for $\det(1 - s_n L_i(u))\sqrt{\det G_i(u)}$, that is,

det
$$D\iota_{\varphi_i}(u, s_n) \neq 0$$
 and det $D\iota_{\varphi_i}(u, s_n) \rightarrow 0$

cf. (2.4). From $D\iota_{\varphi_i}^{-1}(\iota_{\varphi_i}(u,s_n)) = (D\iota_{\varphi_i}(u,s_n))^{-1}$ we conclude

$$\det D\iota_{\varphi_i}^{-1}(\iota_{\varphi_i}(u,s_n)) = \frac{1}{\det D\iota_{\varphi_i}(u,s_n)} \to \infty.$$
(2.7)

On the other hand, by Hypothesis 2.3 (c) the mapping $\iota_{\varphi_i}^{-1}$ is Lipschitz continuous on $\iota_{\varphi_i}(U_i \times (-\beta, \beta))$ and hence $\|D\iota_{\varphi_i}^{-1}\|$ is bounded; this contradicts (2.7).

In the next example we provide a C^2 -hypersurface which does not satisfy Hypothesis 2.3; here a curve in \mathbb{R}^2 with unbounded curvature at "infinity" is discussed.

Example 2.5 Consider the curve

$$\varphi \colon \mathbb{R} \to \mathbb{R}^2, \quad u \mapsto \begin{pmatrix} u \\ \int_0^u \sin(t^2) \mathrm{d}t \end{pmatrix},$$

and observe that $\varphi(\mathbb{R})$ is an orientable C^2 -hypersurface in \mathbb{R}^2 with parametrization $\{\varphi, \mathbb{R}, \mathbb{R}^2\}$. If we fix the unit normal vector field by

$$\nu(u) = \frac{1}{(1 + \sin^2(u^2))^{1/2}} \begin{pmatrix} -\sin(u^2) \\ 1 \end{pmatrix},$$

then the corresponding 1×1 -matrix of the Weingarten map is given by

$$L(u) = \frac{2u\cos(u^2)}{(1+\sin^2(u^2))^{3/2}}$$

Clearly, L is unbounded and hence, item (b) in Hypothesis 2.3 is not satisfied.

Under Hypothesis 2.3, the mapping ι_{Σ} in (2.5) is bijective from $\Sigma \times (-\beta, \beta)$ onto Ω_{β} . This allows us to identify functions f supported on Ω_{β} with functions \tilde{f} defined on $\Sigma \times (-\beta, \beta)$ via the natural identification

$$f(x) = f(\iota_{\Sigma}(x_{\Sigma}, t)) = f(x_{\Sigma}, t), \quad x = \iota_{\Sigma}(x_{\Sigma}, t), \quad x_{\Sigma} \in \Sigma, \ t \in (-\beta, \beta).$$

Subsequently, $L^1(\Omega_\beta)$ is equipped with the *d*-dimensional Lebesgue measure Λ_d and $L^1(\Sigma \times (-\beta, \beta))$ is equipped with the measure $\sigma \times \Lambda_1$. In the next proposition it is shown that $L^1(\Omega_\beta)$ and $L^1(\Sigma \times (-\beta, \beta))$ can be identified and a useful change of variables formula is provided.

Proposition 2.6 Let $\Sigma \subset \mathbb{R}^d$ be an orientable C^2 -hypersurface, assume that Hypothesis 2.3 is satisfied and let Ω_β be as in (2.6). Then the following assertions are true.

(i) Let $\iota_{\beta} := \iota_{\Sigma}|_{\Sigma \times (-\beta,\beta)}$. Then, there exist constants $0 < c_1 \le c_2 < +\infty$ such that

$$c_1 \|f\|_{L^1(\Omega_\beta)} \le \|f \circ \iota_\beta\|_{L^1(\Sigma \times (-\beta,\beta))} \le c_2 \|f\|_{L^1(\Omega_\beta)}, \qquad f \in L^1(\Omega_\beta).$$

In particular, $f \in L^1(\Omega_\beta)$ if and only if $f \circ \iota_\beta \in L^1(\Sigma \times (-\beta, \beta))$.

(ii) For $f \in L^1(\Omega_\beta)$ the identity

$$\int_{\Omega_{\beta}} f(x) dx = \int_{\Sigma} \int_{-\beta}^{\beta} f(x_{\Sigma} + t\nu(x_{\Sigma})) \det(1 - tW(x_{\Sigma})) dt d\sigma(x_{\Sigma})$$

holds, where W is the Weingarten map associated to Σ .

Proof. Let $\{\varphi_i, U_i, V_i\}_{i \in I}$ be a parametrization of Σ with a finite index set I, let $\{\chi_i\}_{i \in I}$ be a partition of unity subordinate to the covering $\{V_i\}_{i \in I}$ and set

$$\widetilde{\chi}_i(x) := \chi_i(x_{\Sigma}), \qquad i \in I,$$

for $x = \iota_{\beta}(x_{\Sigma}, t) = x_{\Sigma} + t\nu(x_{\Sigma}) \in \Omega_{\beta}$ with $x_{\Sigma} \in \Sigma$ and $t \in (-\beta, \beta)$. The family $\{\widetilde{\chi}_i\}_{i \in I}$ satisfies

$$\sum_{i \in I} \widetilde{\chi}_i(x) = 1 \quad \text{for all } x \in \Omega_\beta$$
(2.8)

due to the properties of the partition of unity $\{\chi_i\}_{i \in I}$.

Let $f \in L^1(\Omega_\beta)$, let ι_{φ_i} be as in (2.3) and let $\Omega_{i,\beta} := \iota_{\varphi_i}(U_i \times (-\beta, \beta))$. Using (2.8) we get

$$\int_{\Omega_{\beta}} f(x) \mathrm{d}x = \int_{\Omega_{\beta}} \sum_{i \in I} \widetilde{\chi}_i(x) f(x) \mathrm{d}x = \sum_{i \in I} \int_{\Omega_{i,\beta}} \widetilde{\chi}_i(x) f(x) \mathrm{d}x.$$

Making the substitution $x = \iota_{\varphi_i}(u, t), i \in I$, in each summand of the last formula, we get with the aid of (2.4)

$$\int_{\Omega_{\beta}} f(x) \mathrm{d}x = \sum_{i \in I} \int_{U_i} \int_{-\beta}^{\beta} \widetilde{\chi}_i(\iota_{\varphi_i}(u, t)) f(\iota_{\varphi_i}(u, t)) \det(1 - tL_i(u)) \sqrt{\det G_i(u)} \mathrm{d}t \mathrm{d}u.$$

Using $\tilde{\chi}_i(\iota_{\varphi_i}(u,t)) = \chi_i(\varphi_i(u))$, (2.2) and $\iota_{\varphi_i}(u,t) = \iota_\beta(\varphi_i(u),t)$, we end up with

$$\int_{\Omega_{\beta}} f(x) dx = \int_{-\beta}^{\beta} \int_{\Sigma} f(\iota_{\beta}(x_{\Sigma}, t)) \det(1 - tW(x_{\Sigma})) d\sigma(x_{\Sigma}) dt.$$

Together with Hypothesis 2.3 (b) this formula implies assertions (i) and (ii).

2.2 Schrödinger operators with δ -interactions on hypersurfaces

In this section we recall a representation for the resolvent of the self-adjoint Schrödinger operator $A_{\delta,\alpha}$ in $L^2(\mathbb{R}^d)$ with a δ -interaction supported on the C^2 -hypersurface Σ . As in [10] the operator $A_{\delta,\alpha}$ is defined via the corresponding quadratic form with the help of the first representation theorem [25, Theorem VI 2.1]; the functions in the domain of $A_{\delta,\alpha}$ then satisfy the typical δ -type boundary conditions on Σ , see, e.g., [7, Theorem 3.3]. In the following $A_{\delta,\alpha}$ is the unique self-adjoint operator in $L^2(\mathbb{R}^d)$ associated to the form $\mathfrak{a}_{\delta,\alpha}$ in (2.9); cf. (1.1).

Lemma 2.7 Let $\Sigma \subset \mathbb{R}^d$ be a C^2 -hypersurface, assume that Hypothesis 2.3 is satisfied and let $\alpha \in L^{\infty}(\Sigma)$ be a real valued function. Then the symmetric sesquilinear form

$$\mathfrak{a}_{\delta,\alpha}[f,g] := \left(\nabla f, \nabla g\right)_{L^2(\mathbb{R}^d;\mathbb{C}^d)} - \int_{\Sigma} \alpha f|_{\Sigma} \overline{g|_{\Sigma}} \,\mathrm{d}\sigma, \quad \mathrm{dom}\,\mathfrak{a}_{\delta,\alpha} := H^1(\mathbb{R}^d), \tag{2.9}$$

is densely defined, closed and bounded from below in $L^2(\mathbb{R}^d)$; here $f|_{\Sigma}$, $g|_{\Sigma}$ denote the traces of functions $f, g \in H^1(\mathbb{R}^d)$ on Σ .

Proof. First, we note that Lemma A.3 (i) together with [24, Theorem VII 2, Remark VI 1] (see also [24, Section VIII 1.1]) imply that for $\frac{1}{2} < s \leq 1$ there is a bounded trace operator from $H^s(\mathbb{R}^d)$ to $L^2(\Sigma)$. In particular, the form $\mathfrak{a}_{\delta,\alpha}$ in (2.9) is well defined. Moreover, since $H^1(\mathbb{R}^d)$ is dense in $L^2(\mathbb{R}^d)$, the form $\mathfrak{a}_{\delta,\alpha}$ is densely defined in $L^2(\mathbb{R}^d)$.

Subsequently, fix some $s \in (\frac{1}{2}, 1)$ and $c_s > 0$ such that $||f|_{\Sigma}||_{L^2(\Sigma)} \leq \sqrt{c_s} ||f||_{H^s(\mathbb{R}^d)}$ for all $f \in H^1(\mathbb{R}^d)$. Let $\varepsilon > 0$ and use [23, Theorem 3.30] or [38, Satz 11.18 e)] to see that there exists a $C(\varepsilon) > 0$ such that

$$\begin{aligned} |(\alpha f|_{\Sigma}, f|_{\Sigma})_{L^{2}(\Sigma)}| &\leq \|\alpha\|_{\infty} \|f|_{\Sigma}\|_{L^{2}(\Sigma)}^{2} \\ &\leq c_{s} \|\alpha\|_{\infty} \|f\|_{H^{s}(\mathbb{R}^{d})}^{2} \\ &\leq c_{s} \|\alpha\|_{\infty} \left(\varepsilon \|f\|_{H^{1}(\mathbb{R}^{d})}^{2} + C(\varepsilon)\|f\|_{L^{2}(\mathbb{R}^{d})}^{2}\right). \end{aligned}$$

$$(2.10)$$

Thus, for sufficiently small $\varepsilon > 0$ the form $f \mapsto (\alpha f|_{\Sigma}, f|_{\Sigma})$ on $H^1(\mathbb{R}^d)$ is relatively bounded with respect to the closed and nonnegative form $f \mapsto (\nabla f, \nabla f)_{L^2(\mathbb{R}^d;\mathbb{C}^d)}$ on $H^1(\mathbb{R}^d)$ with bound smaller than one. Then by [25, Theorem VI 1.33] the form $\mathfrak{a}_{\delta,\alpha}$ in (2.9) is closed and bounded from below.

Next, we provide a formula for the resolvent of the Schrödinger operator $A_{\delta,\alpha}$. For this purpose some notations are required. The free Laplace operator in $L^2(\mathbb{R}^d)$ with domain $H^2(\mathbb{R}^d)$ is denoted by $-\Delta$; it is clear that $-\Delta$ coincides with $A_{\delta,0}$ ($\alpha \equiv 0$) in the above lemma. The spectrum of $-\Delta$ is given by $\sigma(-\Delta) = [0, \infty)$. For $\lambda \in (-\infty, 0) \subset \rho(-\Delta)$ we define the function

$$G_{\lambda}(x) := \frac{1}{(2\pi)^{d/2}} \left(\frac{|x|}{\sqrt{-\lambda}}\right)^{1-d/2} K_{d/2-1}\left(\sqrt{-\lambda}|x|\right), \quad x \in \mathbb{R}^d \setminus \{0\},$$
(2.11)

where $K_{d/2-1}$ denotes a modified Bessel function of the second kind and order $\frac{d}{2} - 1$, see [1] for the definition and the properties of these functions. Then,

$$(R(\lambda)f)(x) := \left((-\Delta - \lambda)^{-1}f\right)(x) = \int_{\mathbb{R}^d} G_\lambda(x - y)f(y) \mathrm{d}y;$$

cf. [36, Section 7.4]. Next we define for $\lambda \in (-\infty, 0)$ integral operators $\gamma(\lambda)$, $M(\lambda)$ and provide the integral representation for the adjoint of $\gamma(\lambda)$

$$\gamma(\lambda): L^2(\Sigma) \to L^2(\mathbb{R}^d), \qquad (\gamma(\lambda)\xi)(x):=\int_{\Sigma} G_\lambda(x-y_{\Sigma})\xi(y_{\Sigma})\mathrm{d}\sigma(y_{\Sigma}); \qquad (2.12a)$$

$$M(\lambda): L^{2}(\Sigma) \to L^{2}(\Sigma), \qquad (M(\lambda)\xi)(x_{\Sigma}):= \int_{\Sigma} G_{\lambda}(x_{\Sigma} - y_{\Sigma})\xi(y_{\Sigma}) \mathrm{d}\sigma(y_{\Sigma}); \qquad (2.12b)$$

$$\gamma(\lambda)^* : L^2(\mathbb{R}^d) \to L^2(\Sigma), \qquad (\gamma(\lambda)^* f)(x_{\Sigma}) = \int_{\mathbb{R}^d} G_\lambda(x_{\Sigma} - y) f(y) \mathrm{d}y. \tag{2.12c}$$

For our later considerations the resolvent formula in the next theorem is particularly useful. In the proof of item (a) and later in Section 3 the Schur test for integral operators will be used frequently, see, e.g., [25, Example III 2.4] or [38, Satz 6.9].

Theorem 2.8 Let $\Sigma \subset \mathbb{R}^d$ be a C^2 -hypersurface which satisfies Hypothesis 2.3 and let $\alpha \in L^{\infty}(\Sigma)$ be a real valued function. Then the following statements are true.

- (a) For $\lambda \in (-\infty, 0)$ the operators $\gamma(\lambda), M(\lambda)$ and $\gamma(\lambda)^*$ in (2.12) are bounded and everywhere defined.
- (b) There exists a λ₀ < 0 such that 1 − αM(λ) admits a bounded and everywhere defined inverse for all λ ∈ (−∞, λ₀). These λ belong to ρ(A_{δ,α}) and it holds

$$(A_{\delta,\alpha} - \lambda)^{-1} = R(\lambda) + \gamma(\lambda) (1 - \alpha M(\lambda))^{-1} \alpha \gamma(\lambda)^*.$$

Proof. (a) Let $\lambda \in (-\infty, 0)$. In order to prove that $\gamma(\lambda)$ is well defined and bounded we use the Schur test. In fact, from Proposition A.4 (i) and Proposition A.2 (i) we obtain

$$\|\gamma(\lambda)\|^2 \leq \sup_{x \in \mathbb{R}^d} \int_{\Sigma} \left| G_{\lambda}(x - y_{\Sigma}) \right| \mathrm{d}\sigma(y_{\Sigma}) \cdot \sup_{y_{\Sigma} \in \Sigma} \int_{\mathbb{R}^d} \left| G_{\lambda}(x - y_{\Sigma}) \right| \mathrm{d}x < \infty.$$

In a similar way one can show that $M(\lambda)$ and $\gamma(\lambda)^*$ are bounded.

Item (b) is essentially a variant of [10, Lemma 2.3]. In fact, let us define for Borel sets $B \subset \mathbb{R}^d$ the measure m by

$$m(B) := \sigma(B \cap \Sigma), \tag{2.13}$$

where σ is the measure in (2.2). Then $m(\mathbb{R}^d \setminus \Sigma) = 0$, the spaces $L^2(\mathbb{R}^d; m)$ and $L^2(\Sigma)$ can be identified and for $f \in L^1(\Sigma)$ one has

$$\int_{\Sigma} f(x_{\Sigma}) \mathrm{d}\sigma(x_{\Sigma}) = \int_{\mathbb{R}^d} \widetilde{f}(x) \, \mathrm{d}m(x)$$

where \tilde{f} is some extension of f onto the m-null set $\mathbb{R}^d \setminus \Sigma$. Moreover, the estimate (2.10) shows that the measure m in (2.13) satisfies [10, eq. (2.1)] (with $\gamma = -\alpha$). Now it is easy to see that the integral operators $\gamma(\lambda)$, $M(\lambda)$, and $\gamma(\lambda)^*$ in (2.12) can be identified with the operators $R_{m dx}(i\sqrt{-\lambda})$, $R_{mm}(i\sqrt{-\lambda})$ and $R_{dxm}(i\sqrt{-\lambda})$ in [10], respectively. The assertion in item (b) follows from [10, Lemma 2.3 (ii) and (iii)].

3 Approximation of $A_{\delta,\alpha}$ by Schrödinger operators with regular potentials

In this section we prove Theorem 1.1, the main result of this paper. First, in Section 3.1 we recall briefly the definitions of the layer neighborhoods, the scaled potentials and the associated Hamiltonians H_{ε} from the introduction, and we derive a resolvent formula for H_{ε} which is convenient in the convergence analysis. Section 3.2 contains the main part of the proof of Theorem 1.1. It is efficient to prove Theorem 1.1 for the special case Q = 0 first; all technical estimates in Appendix A and all preparatory steps in Sections 3.1 and 3.2 are tailormade for this case. The general case $Q \neq 0$ is treated with a simple perturbation argument in the last step of the proof of Theorem 1.1. Finally, in Section 3.3 we discuss some connections between the spectral properties of H_{ε} and of $A_{\delta,\alpha}$ that follow from Theorem 1.1.

3.1 Preliminary considerations on H_{ε}

Let $d \ge 2$ and $\Sigma \subset \mathbb{R}^d$ be a C^2 -hypersurface which satisfies Hypothesis 2.3. Then the mapping ι_{Σ} in (2.5) is injective on $\Sigma \times (-\beta, \beta)$ for some (in the following fixed) $\beta > 0$ as in Hypothesis 2.3. Recall the definition of the layer Ω_{ε} from (2.6) for $\varepsilon \in (0, \beta]$, fix a real valued potential $V \in L^{\infty}(\mathbb{R}^d)$ with supp $V \subset \Omega_{\beta}$ and consider the scaled potentials

$$V_{\varepsilon}(x) = \begin{cases} \frac{\beta}{\varepsilon} V\left(x_{\Sigma} + \frac{\beta}{\varepsilon} t\nu(x_{\Sigma})\right), & \text{if } x = x_{\Sigma} + t\nu(x_{\Sigma}) \in \Omega_{\varepsilon}, \\ 0, & \text{else,} \end{cases}$$

where ν is the continuous unit normal vector field on Σ . Observe that $V_{\beta} = V$ and supp $V_{\varepsilon} \subset \Omega_{\varepsilon}$. The associated self-adjoint Schrödinger operators are given by

$$H_{\varepsilon}f = -\Delta f - V_{\varepsilon}f, \quad \operatorname{dom} H_{\varepsilon} = H^2(\mathbb{R}^d).$$
(3.1)

Our main objective in this section is to derive the resolvent formula for H_{ε} in Proposition 3.3 which turns out to be particularly convenient for our convergence analysis. We start with the standard factorization of the potentials $V_{\varepsilon} = v_{\varepsilon} u_{\varepsilon}$, where

$$u_{\varepsilon}: L^{2}(\mathbb{R}^{d}) \to L^{2}(\Omega_{\varepsilon}), \quad (u_{\varepsilon}f)(x) := |V_{\varepsilon}(x)|^{1/2}f(x), \quad x \in \Omega_{\varepsilon},$$
(3.2)

and

$$v_{\varepsilon}: L^{2}(\Omega_{\varepsilon}) \to L^{2}(\mathbb{R}^{d}), \quad (v_{\varepsilon}h)(x) := \begin{cases} \operatorname{sign} V_{\varepsilon}(x)|V_{\varepsilon}(x)|^{1/2}h(x), & x \in \Omega_{\varepsilon}, \\ 0, & \text{else.} \end{cases}$$
(3.3)

Recall that for $\lambda \in \rho(-\Delta) = \mathbb{C} \setminus [0, \infty)$ the resolvent of $-\Delta$ is denoted by $R(\lambda) = (-\Delta - \lambda)^{-1}$. The following proposition contains a first auxiliary resolvent formula for H_{ε} .

Proposition 3.1 Let H_{ε} be defined as in (3.1) and let $u_{\varepsilon}, v_{\varepsilon}$ and $R(\lambda)$ be given as above. Then the following assertions are true.

(i) For all $\lambda \in \mathbb{C} \setminus [0, \infty)$ with $1 \in \rho(u_{\varepsilon}R(\lambda)v_{\varepsilon})$ one has $\lambda \in \rho(H_{\varepsilon})$ and

$$(H_{\varepsilon} - \lambda)^{-1} = R(\lambda) + R(\lambda)v_{\varepsilon} (1 - u_{\varepsilon}R(\lambda)v_{\varepsilon})^{-1} u_{\varepsilon}R(\lambda).$$

(ii) For all $M \in (0,1)$ there exists $\lambda_M < 0$ such that

$$\|u_{\varepsilon}R(\lambda)v_{\varepsilon}\| \le M$$

holds for all $\varepsilon \in (0,\beta]$ and $\lambda < \lambda_M$. In particular, for these λ the results from (i) apply and hence $(-\infty, \lambda_M) \subset \rho(H_{\varepsilon})$ for all $\varepsilon \in (0,\beta]$.

Proof. (i) Let $\lambda \in \mathbb{C} \setminus [0, \infty)$ be such that $1 \in \rho(u_{\varepsilon}R(\lambda)v_{\varepsilon})$. Note that λ is not an eigenvalue of H_{ε} , as otherwise $1 \in \sigma_p(u_{\varepsilon}R(\lambda)v_{\varepsilon})$; cf. [9, Lemma 1]. Next, we define the operator

$$T(\lambda) := R(\lambda) + R(\lambda)v_{\varepsilon} \left(1 - u_{\varepsilon}R(\lambda)v_{\varepsilon}\right)^{-1} u_{\varepsilon}R(\lambda).$$

This operator is well defined and bounded, as $1 \in \rho(u_{\varepsilon}R(\lambda)v_{\varepsilon})$. Using $V_{\varepsilon} = v_{\varepsilon}u_{\varepsilon}$ we conclude that

$$(H_{\varepsilon} - \lambda)T(\lambda)f = (-\Delta - \lambda - v_{\varepsilon}u_{\varepsilon})T(\lambda)f$$

= $f + v_{\varepsilon} (1 - u_{\varepsilon}R(\lambda)v_{\varepsilon})^{-1} u_{\varepsilon}R(\lambda)f - v_{\varepsilon}u_{\varepsilon}R(\lambda)f$
 $- v_{\varepsilon}(1 - 1 + u_{\varepsilon}R(\lambda)v_{\varepsilon}) (1 - u_{\varepsilon}R(\lambda)v_{\varepsilon})^{-1} u_{\varepsilon}R(\lambda)f$
= $f + v_{\varepsilon} (1 - u_{\varepsilon}R(\lambda)v_{\varepsilon})^{-1} u_{\varepsilon}R(\lambda)f - v_{\varepsilon}u_{\varepsilon}R(\lambda)f$
 $- v_{\varepsilon} (1 - u_{\varepsilon}R(\lambda)v_{\varepsilon})^{-1} u_{\varepsilon}R(\lambda)f + v_{\varepsilon}u_{\varepsilon}R(\lambda)f$
= f

holds for any $f \in L^2(\mathbb{R}^d)$. Hence, $(H_{\varepsilon} - \lambda)$ is bijective, which implies that $\lambda \in \rho(H_{\varepsilon})$, and

$$(H_{\varepsilon} - \lambda)^{-1} = T(\lambda) = R(\lambda) + R(\lambda)v_{\varepsilon} \left(1 - u_{\varepsilon}R(\lambda)v_{\varepsilon}\right)^{-1} u_{\varepsilon}R(\lambda).$$

This proves assertion (i).

(ii) Let $\lambda \in (-\infty, 0)$ and recall that $R(\lambda)$ can be expressed by $R(\lambda)f = \int_{\mathbb{R}^d} G_{\lambda}(\cdot - y)f(y) dy$ with G_{λ} as in (2.11). Let $M \in (0,1)$ be fixed. Using the Schur test and that the absolute value of the integral kernel of $u_{\varepsilon}R(\lambda)v_{\varepsilon}$ is symmetric, we find that

$$\left\| u_{\varepsilon} R(\lambda) v_{\varepsilon} \right\| \leq \sup_{x \in \Omega_{\varepsilon}} \int_{\Omega_{\varepsilon}} |V_{\varepsilon}(x)|^{1/2} |G_{\lambda}(x-y)| |V_{\varepsilon}(y)|^{1/2} \mathrm{d}y \leq \frac{\beta}{\varepsilon} \|V\|_{L^{\infty}} \sup_{x \in \Omega_{\varepsilon}} \int_{\Omega_{\varepsilon}} |G_{\lambda}(x-y)| \mathrm{d}y.$$

Hence, the claimed result follows from Proposition A.4 (ii), as this shows the existence of a number $\lambda_M < 0$ such that

$$\frac{\beta}{\varepsilon} \|V\|_{L^{\infty}} \int_{\Omega_{\varepsilon}} |G_{\lambda}(x-y)| \mathrm{d}y \le M$$

for any $x \in \Omega_{\varepsilon}$, all $\lambda < \lambda_M$ and any $\varepsilon \in (0, \beta]$.

Next, we transform the resolvent formula from Proposition 3.1 into another one, which is more convenient for the convergence analysis. This requires several preparatory steps. Recall that for an interval $I \subset \mathbb{R}$ the space $L^2(\Sigma \times I)$ is equipped with the product measure $\sigma \times \Lambda_1$ of the Hausdorff measure on Σ and the one-dimensional Lebesgue measure. Define the functions $u, v \in L^{\infty}(\Sigma \times (-1, 1))$ by

$$u(x_{\Sigma},t) := |\beta V(x_{\Sigma} + \beta t \nu(x_{\Sigma}))|^{1/2} \quad \text{and} \quad v(x_{\Sigma},t) := \operatorname{sign}\left(V(x_{\Sigma} + \beta t \nu(x_{\Sigma}))\right) \cdot u(x_{\Sigma},t).$$
(3.4)

The following operators are essential to state a convenient resolvent formula for H_{ε} . We define for $\varepsilon \in [0,\beta]$ and $\lambda \in (-\infty, 0)$ the integral operators $A_{\varepsilon}(\lambda) \colon L^2(\Sigma \times (-1, 1)) \to L^2(\mathbb{R}^d), B_{\varepsilon}(\lambda) \colon L^2(\Sigma \times (-1, 1)) \to L^2(\Sigma \times (-1, 1))$ and $C_{\varepsilon}(\lambda) \colon L^2(\mathbb{R}^d) \to L^2(\Sigma \times (-1, 1))$ as

$$(A_{\varepsilon}(\lambda)\Xi)(x) := \int_{\Sigma} \int_{-1}^{1} G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma}))v(y_{\Sigma}, s)$$

$$\cdot \det(1 - \varepsilon sW(y_{\Sigma}))\Xi(y_{\Sigma}, s)dsd\sigma(y_{\Sigma});$$

$$(B_{\varepsilon}(\lambda)\Xi)(x_{\Sigma}, t) := u(x_{\Sigma}, t) \int_{\Sigma} \int_{-1}^{1} G_{\lambda}(x_{\Sigma} + \varepsilon t\nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma}))$$

$$\cdot v(y_{\Sigma}, s) \det(1 - \varepsilon sW(y_{\Sigma}))\Xi(y_{\Sigma}, s)dsd\sigma(y_{\Sigma});$$

$$(C_{\varepsilon}(\lambda)f)(x_{\Sigma}, t) := u(x_{\Sigma}, t) \int_{\mathbb{R}^{d}} G_{\lambda}(x_{\Sigma} + \varepsilon t\nu(x_{\Sigma}) - y)f(y)dy.$$

$$(3.5a)$$

In order to investigate the properties of
$$A_{\varepsilon}(\lambda), B_{\varepsilon}(\lambda)$$
 and $C_{\varepsilon}(\lambda)$, we introduce several auxiliary operators. For $\varepsilon \in (0, \beta]$ define the embedding operator

$$\mathcal{I}_{\varepsilon}: L^{2}(\Sigma \times (-\varepsilon, \varepsilon)) \to L^{2}(\Omega_{\varepsilon}), \quad (\mathcal{I}_{\varepsilon}\Phi)(x_{\Sigma} + t\nu(x_{\Sigma})) := \Phi(x_{\Sigma}, t).$$
(3.6)

It follows from Hypothesis 2.3 (b) and Proposition 2.6 that the operator $\mathcal{I}_{\varepsilon}$ is bounded, everywhere defined and bijective. Its inverse is given by

$$\mathcal{I}_{\varepsilon}^{-1}: L^{2}(\Omega_{\varepsilon}) \to L^{2}(\Sigma \times (-\varepsilon, \varepsilon)), \quad (\mathcal{I}_{\varepsilon}^{-1}h)(x_{\Sigma}, t) = h(x_{\Sigma} + t\nu(x_{\Sigma})).$$

Furthermore, we consider the scaling operator

In

$$S_{\varepsilon}: L^{2}(\Sigma \times (-1,1)) \to L^{2}(\Sigma \times (-\varepsilon,\varepsilon)), \quad (S_{\varepsilon}\Xi)(x_{\Sigma},t) := \frac{1}{\sqrt{\varepsilon}}\Xi\left(x_{\Sigma},\frac{t}{\varepsilon}\right).$$
(3.7)

The operator S_{ε} is unitary and its inverse is

$$S_{\varepsilon}^{-1}: L^{2}(\Sigma \times (-\varepsilon, \varepsilon)) \to L^{2}(\Sigma \times (-1, 1)), \qquad (S_{\varepsilon}^{-1}\Phi)(x_{\Sigma}, t) = \sqrt{\varepsilon}\Phi(x_{\Sigma}, \varepsilon t).$$

In the next lemma some properties of $A_{\varepsilon}(\lambda), B_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ are provided.

Lemma 3.2 Let $\lambda \in (-\infty, 0)$ and let $\varepsilon \in (0, \beta]$. Moreover, let the operators $u_{\varepsilon}, v_{\varepsilon}, \mathcal{I}_{\varepsilon}$ and S_{ε} be defined by (3.2), (3.3), (3.6) and (3.7), respectively, and let the integral operators $A_{\varepsilon}(\lambda)$, $B_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ be as in (3.5). Then the following assertions are true.

(i) It holds

$$A_{\varepsilon}(\lambda) = R(\lambda)v_{\varepsilon}\mathcal{I}_{\varepsilon}S_{\varepsilon} \quad and \quad C_{\varepsilon}(\lambda) = S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1}u_{\varepsilon}R(\lambda).$$

In particular, the operators $A_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ are bounded and everywhere defined.

(ii) It holds

$$B_{\varepsilon}(\lambda) = S_{\varepsilon}^{-1} \mathcal{I}_{\varepsilon}^{-1} u_{\varepsilon} R(\lambda) v_{\varepsilon} \mathcal{I}_{\varepsilon} S_{\varepsilon}$$

Moreover, for any $M \in (0, 1)$ there exists $\lambda_M < 0$ such that $||B_{\varepsilon}(\lambda)|| \le M$ for all $\lambda \in (-\infty, \lambda_M)$ and all $\varepsilon \in (0, \beta]$. In particular, $B_{\varepsilon}(\lambda)$ is bounded and everywhere defined and for $\lambda < \lambda_M$ the operator $1 - B_{\varepsilon}(\lambda)$ has a bounded and everywhere defined inverse.

Proof. Let $\varepsilon \in (0, \beta]$ and $\lambda \in (-\infty, 0)$ be fixed.

(i) We show the two formulae $A_{\varepsilon}(\lambda) = R(\lambda)v_{\varepsilon}\mathcal{I}_{\varepsilon}S_{\varepsilon}$ and $C_{\varepsilon}(\lambda) = S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1}u_{\varepsilon}R(\lambda)$. The operators $A_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ are then automatically bounded and everywhere defined, as the operators S_{ε} , $\mathcal{I}_{\varepsilon}$ as well as their inverses, u_{ε} , v_{ε} and $R(\lambda)$ have these properties.

Let $\Xi \in L^2(\Sigma \times (-1, 1))$. By the definitions of $\mathcal{I}_{\varepsilon}$ and S_{ε} it holds that

$$(\mathcal{I}_{\varepsilon}S_{\varepsilon}\Xi)(y_{\Sigma}+s\nu(y_{\Sigma}))=\frac{1}{\sqrt{\varepsilon}}\Xi\left(y_{\Sigma},\frac{s}{\varepsilon}\right)$$

Furthermore, the definitions of v_{ε} and v (see (3.4)) imply for any $h \in L^2(\Omega_{\varepsilon})$ that

$$(v_{\varepsilon}h) (y_{\Sigma} + s\nu(y_{\Sigma})) = \frac{1}{\sqrt{\varepsilon}} \operatorname{sign} V \left(y_{\Sigma} + \frac{\beta}{\varepsilon} s\nu(y_{\Sigma}) \right) \left| \beta V \left(y_{\Sigma} + \frac{\beta}{\varepsilon} s\nu(y_{\Sigma}) \right) \right|^{1/2} h \left(y_{\Sigma} + s\nu(y_{\Sigma}) \right)$$
$$= \frac{1}{\sqrt{\varepsilon}} v \left(y_{\Sigma}, \frac{s}{\varepsilon} \right) h \left(y_{\Sigma} + s\nu(y_{\Sigma}) \right)$$

for a.e. $(y_{\Sigma}, s) \in \Sigma \times (-\varepsilon, \varepsilon)$. Using the transformation $\Omega_{\varepsilon} \ni y = y_{\Sigma} + s\nu(y_{\Sigma}) \mapsto (y_{\Sigma}, s) \in \Sigma \times (-\varepsilon, \varepsilon)$ and Proposition 2.6 we find that

$$(R(\lambda)v_{\varepsilon}\mathcal{I}_{\varepsilon}S_{\varepsilon}\Xi)(x) = \int_{\mathbb{R}^{d}} G_{\lambda}(x-y)(v_{\varepsilon}\mathcal{I}_{\varepsilon}S_{\varepsilon}\Xi)(y)dy = \int_{\Sigma} \int_{-\varepsilon}^{\varepsilon} G_{\lambda}(x-y_{\Sigma}-s\nu(y_{\Sigma}))\frac{1}{\varepsilon}(v\Xi)\left(y_{\Sigma},\frac{s}{\varepsilon}\right)\det(1-sW(y_{\Sigma}))dsd\sigma(y_{\Sigma}) = \int_{\Sigma} \int_{-1}^{1} G_{\lambda}(x-y_{\Sigma}-\varepsilon r\nu(y_{\Sigma}))v\left(y_{\Sigma},r\right)\det(1-\varepsilon rW(y_{\Sigma}))\Xi\left(y_{\Sigma},r\right)drd\sigma(y_{\Sigma}) = (A_{\varepsilon}(\lambda)\Xi)(x).$$

$$(3.8)$$

Since this is true for a.e. $x \in \mathbb{R}^d$, the first formula of item (i) is shown.

Next, we show the assertion on $C_{\varepsilon}(\lambda)$. A simple calculation yields that

$$(S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1}u_{\varepsilon}g)(x_{\Sigma},t) = \sqrt{\varepsilon}(\mathcal{I}_{\varepsilon}^{-1}u_{\varepsilon}g)(x_{\Sigma},\varepsilon t)$$

$$= \sqrt{\varepsilon} \cdot \frac{1}{\sqrt{\varepsilon}} \left| \beta V \left(x_{\Sigma} + \frac{\beta}{\varepsilon}\varepsilon t\nu(x_{\Sigma}) \right) \right|^{1/2} g(x_{\Sigma} + \varepsilon t\nu(x_{\Sigma}))$$

$$= u(x_{\Sigma},t)g(x_{\Sigma} + \varepsilon t\nu(x_{\Sigma}))$$
(3.9)

for any $g \in L^2(\mathbb{R}^d)$ and a.e. $(x_{\Sigma}, t) \in \Sigma \times (-1, 1)$. Using (3.9), we find that

$$(S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1}u_{\varepsilon}R(\lambda)f)(x_{\Sigma},t) = u(x_{\Sigma},t)\int_{\mathbb{R}^{d}}G_{\lambda}(x_{\Sigma}+\varepsilon t\nu(x_{\Sigma})-y)f(y)\mathrm{d}y = (C_{\varepsilon}(\lambda)f)(x_{\Sigma},t)$$

for all $f \in L^2(\mathbb{R}^d)$ and a.e. $(x_{\Sigma}, t) \in \Sigma \times (-1, 1)$. Thus, the second formula is shown as well.

(ii) Using (3.9) and (3.8) we get that

$$\begin{split} \left(S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1}u_{\varepsilon}R(\lambda)v_{\varepsilon}\mathcal{I}_{\varepsilon}S_{\varepsilon}\Xi\right)(x_{\Sigma},t) &= u(x_{\Sigma},t)\left(R(\lambda)v_{\varepsilon}\mathcal{I}_{\varepsilon}S_{\varepsilon}\Xi\right)(x_{\Sigma}+\varepsilon t\nu(x_{\Sigma}))\\ &= u(x_{\Sigma},t)\int_{\Sigma}\int_{-1}^{1}G_{\lambda}(x_{\Sigma}+\varepsilon t\nu(x_{\Sigma})-y_{\Sigma}-\varepsilon r\nu(y_{\Sigma}))\\ &\quad \cdot v\left(y_{\Sigma},r\right)\det(1-\varepsilon rW(y_{\Sigma}))\Xi\left(y_{\Sigma},r\right)\mathrm{d}r\mathrm{d}\sigma(y_{\Sigma})\\ &= (B_{\varepsilon}(\lambda)\Xi)(x_{\Sigma},t). \end{split}$$

Therefore, we obtain the desired formula in item (ii). In particular, this formula implies that $B_{\varepsilon}(\lambda)$ is bounded and everywhere defined.

Let $M \in (0,1)$ be fixed. Note that $\mathcal{I}_{\varepsilon}$ and $\mathcal{I}_{\varepsilon}^{-1}$ are uniformly bounded for all $\varepsilon \in (0,\beta]$ by Proposition 2.6 and Hypothesis 2.3 (b). Furthermore, recall that S_{ε} is unitary. Hence, using Proposition 3.1 (ii) we find that

$$\|B_{\varepsilon}(\lambda)\| = \|S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1}u_{\varepsilon}R(\lambda)v_{\varepsilon}\mathcal{I}_{\varepsilon}S_{\varepsilon}\| \le \|S_{\varepsilon}^{-1}\| \cdot \|\mathcal{I}_{\varepsilon}^{-1}\| \cdot \|u_{\varepsilon}R(\lambda)v_{\varepsilon}\| \cdot \|\mathcal{I}_{\varepsilon}\| \cdot \|S_{\varepsilon}\| \le M$$

for all $\lambda < 0$ with $|\lambda|$ sufficiently large and all $\varepsilon \in (0, \beta]$.

After all these preparatory steps it is simple to transform the resolvent formula for H_{ε} from Proposition 3.1 into another one which is more convenient for the convergence analysis.

Proposition 3.3 Let H_{ε} be defined as in (3.1) and let $A_{\varepsilon}(\lambda), B_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ be as in (3.5). Then there exists a $\lambda_0 < 0$ such that $(-\infty, \lambda_0) \subset \rho(H_{\varepsilon})$ and

$$(H_{\varepsilon} - \lambda)^{-1} = R(\lambda) + A_{\varepsilon}(\lambda) (1 - B_{\varepsilon}(\lambda))^{-1} C_{\varepsilon}(\lambda)$$

for all $\varepsilon \in (0, \beta]$ and $\lambda < \lambda_0$.

Proof. Let the operators u_{ε} , v_{ε} , $\mathcal{I}_{\varepsilon}$, S_{ε} be defined as in (3.2), (3.3), (3.6) and (3.7), respectively. Choose $\lambda_0 < 0$ such that $1 - u_{\varepsilon}R(\lambda)v_{\varepsilon}$ and $1 - B_{\varepsilon}(\lambda)$ are boundedly invertible for any $\lambda < \lambda_0$ and all $\varepsilon \in (0, \beta]$ (such a λ_0 exists by Proposition 3.1 and Lemma 3.2). Then, it holds by Proposition 3.1 and Lemma 3.2 that

$$(H_{\varepsilon} - \lambda)^{-1} = R(\lambda) + R(\lambda)v_{\varepsilon} (1 - u_{\varepsilon}R(\lambda)v_{\varepsilon})^{-1} u_{\varepsilon}R(\lambda)$$

= $R(\lambda) + A_{\varepsilon}(\lambda)S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1} (1 - \mathcal{I}_{\varepsilon}S_{\varepsilon}B_{\varepsilon}(\lambda)S_{\varepsilon}^{-1}\mathcal{I}_{\varepsilon}^{-1})^{-1}\mathcal{I}_{\varepsilon}S_{\varepsilon}C_{\varepsilon}(\lambda)$
= $R(\lambda) + A_{\varepsilon}(\lambda) (1 - B_{\varepsilon}(\lambda))^{-1} C_{\varepsilon}(\lambda),$

which proves the statement of this proposition.

3.2 Proof of Theorem 1.1

The main idea to prove Theorem 1.1 is to consider first the special case Q = 0. In this situation we have to show that the family of operators H_{ε} converges in the norm resolvent sense to $A_{\delta,\alpha}$, as $\varepsilon \to 0+$. Because of Theorem 2.8 and Proposition 3.3 it is sufficient to investigate the convergence of $A_{\varepsilon}(\lambda), B_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ separately. This is done in the following lemma. In the proof we make use of the integral estimates in Appendix A and we frequently use the Schur test for integral operators; cf. [25, Example III 2.4] or [38, Satz 6.9]. The general case $Q \neq 0$ follows then by a simple perturbation argument.

Lemma 3.4 Let $\lambda \in (-\infty, 0)$ and let $A_{\varepsilon}(\lambda)$, $B_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ be defined as in (3.5). Then there exists a constant $c = c(d, \lambda, \Sigma, V) > 0$ such that for all sufficiently small $\varepsilon > 0$ the following estimates hold:

$$\|A_{\varepsilon}(\lambda) - A_{0}(\lambda)\|, \|C_{\varepsilon}(\lambda) - C_{0}(\lambda)\| \le c\varepsilon (1 + |\ln \varepsilon|)^{1/2}, \qquad \|B_{\varepsilon}(\lambda) - B_{0}(\lambda)\| \le c\varepsilon (1 + |\ln \varepsilon|).$$

Proof. First, we provide an estimate related to the Weingarten map $W(y_{\Sigma})$. Let $\mu_1(y_{\Sigma}), \ldots, \mu_{d-1}(y_{\Sigma})$ be the eigenvalues of the matrix of $W(y_{\Sigma})$, which are independent of the parametrization of Σ and uniformly bounded on Σ ; cf. Proposition 2.4. This implies for $y_{\Sigma} \in \Sigma$, $s \in (-1, 1)$ and $\varepsilon \in (0, 1)$ the existence of a constant $c_1 > 0$ such that

$$D_{\varepsilon}(y_{\Sigma}, s) := \left|1 - \det(1 - \varepsilon s W(y_{\Sigma}))\right| = \left|1 - \prod_{k=1}^{d-1} (1 - \varepsilon s \mu_k(y_{\Sigma}))\right| \le c_1 \varepsilon.$$
(3.10)

Fix $\lambda \in (-\infty, 0)$. In order to find an estimate for $||A_{\varepsilon}(\lambda) - A_0(\lambda)||$, we introduce for $\varepsilon > 0$ the auxiliary integral operator $\widehat{A}_{\varepsilon}(\lambda)$: $L^2(\Sigma \times (-1, 1)) \to L^2(\mathbb{R}^d)$ by

$$\left(\widehat{A}_{\varepsilon}(\lambda)\Xi\right)(x) := \int_{\Sigma} \int_{-1}^{1} G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma}))v(y_{\Sigma}, s)\Xi(y_{\Sigma}, s)\mathrm{d}s\mathrm{d}\sigma(y_{\Sigma}).$$

The quantities $||A_{\varepsilon}(\lambda) - \widehat{A}_{\varepsilon}(\lambda)||$ and $||\widehat{A}_{\varepsilon}(\lambda) - A_0(\lambda)||$ are estimated separately for sufficiently small $\varepsilon > 0$. To estimate $||\widehat{A}_{\varepsilon}(\lambda) - A_{\varepsilon}(\lambda)||$ we introduce the functions $\widetilde{F}_{1,\varepsilon}^A : \mathbb{R}^d \to [0,\infty]$ and $\widehat{F}_{1,\varepsilon}^A : \Sigma \times (-1,1) \to [0,\infty]$ by

$$\widetilde{F}_{1,\varepsilon}^{A}(x) := \int_{\Sigma} \int_{-1}^{1} \left| G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma}))v(y_{\Sigma}, s) \right| D_{\varepsilon}(y_{\Sigma}, s) \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}),$$
$$\widehat{F}_{1,\varepsilon}^{A}(y_{\Sigma}, s) := \int_{\mathbb{R}^{d}} \left| G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma}))v(y_{\Sigma}, s) \right| D_{\varepsilon}(y_{\Sigma}, s) \mathrm{d}x.$$

With the aid of (3.10), Proposition A.5 (i) and Proposition A.2 (ii), we find that

$$\sup \widetilde{F}_{1,\varepsilon}^A \le c_{A,1}\varepsilon \quad \text{and} \quad \sup \widehat{F}_{1,\varepsilon}^A \le c_{A,1}\varepsilon$$

with a constant $c_{A,1} = c_{A,1}(d, \lambda, \Sigma, V) > 0$. Using these bounds and the Schur test we obtain

$$\left\|\widehat{A}_{\varepsilon}(\lambda) - A_{\varepsilon}(\lambda)\right\|^{2} \le \left(\sup \widetilde{F}_{1,\varepsilon}^{A}\right) \cdot \left(\sup \widehat{F}_{1,\varepsilon}^{A}\right) \le c_{A,1}^{2}\varepsilon^{2}.$$
(3.11)

To estimate $\|\hat{A}_{\varepsilon}(\lambda) - A_0(\lambda)\|$, we introduce the functions $\widetilde{F}_{2,\varepsilon}^A \colon \mathbb{R}^d \to [0,\infty]$ and $\widehat{F}_{2,\varepsilon}^A \colon \Sigma \times (-1,1) \to [0,\infty]$ by

$$\widetilde{F}_{2,\varepsilon}^{A}(x) := \int_{\Sigma} \int_{-1}^{1} \left| \left(G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma})) - G_{\lambda}(x - y_{\Sigma}) \right) v(y_{\Sigma}, s) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma})$$
$$\widehat{F}_{2,\varepsilon}^{A}(y_{\Sigma}, s) := \int_{\mathbb{R}^{d}} \left| \left(G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma})) - G_{\lambda}(x - y_{\Sigma}) \right) v(y_{\Sigma}, s) \right| \mathrm{d}s.$$

We find with the help of Proposition A.9 (i) and Proposition A.6 that

$$\sup \widetilde{F}_{2,\varepsilon}^{A} \leq c_{A,2}\varepsilon(1+|\ln\varepsilon|) \quad \text{and} \quad \sup \widehat{F}_{2,\varepsilon}^{A} \leq c_{A,2}\varepsilon$$

with a constant $c_{A,2} = c_{A,2}(d, \lambda, \Sigma, V) > 0$. Using these bounds and the Schur test we get

$$\left\|\widehat{A}_{\varepsilon}(\lambda) - A_0(\lambda)\right\|^2 \le \left(\sup \widetilde{F}_{2,\varepsilon}^A\right) \cdot \left(\sup \widehat{F}_{2,\varepsilon}^A\right) \le c_{A,2}^2 \varepsilon^2 (1 + |\ln \varepsilon|).$$
(3.12)

Combining the estimates (3.11), (3.12) and applying the triangle inequality for the operator norm, we conclude that there exists a constant $c_A = c_A(d, \lambda, \Sigma, V) > 0$ such that

$$\begin{aligned} \|A_{\varepsilon}(\lambda) - A_{0}(\lambda)\| &\leq \|A_{\varepsilon}(\lambda) - \widehat{A}_{\varepsilon}(\lambda)\| + \|\widehat{A}_{\varepsilon}(\lambda) - A_{0}(\lambda)\| \\ &\leq c_{A,1}\varepsilon + c_{A,2}\varepsilon(1 + |\ln\varepsilon|)^{1/2} \\ &\leq c_{A}\varepsilon(1 + |\ln\varepsilon|)^{1/2}. \end{aligned}$$

Next, we analyze the convergence of $B_{\varepsilon}(\lambda)$. For this purpose, we introduce for $\varepsilon > 0$ the auxiliary operator $\widehat{B}_{\varepsilon}(\lambda): L^2(\Sigma \times (-1,1)) \to L^2(\Sigma \times (-1,1))$ as

$$\left(\widehat{B}_{\varepsilon}(\lambda)\Xi\right)(x_{\Sigma},t) := u(x_{\Sigma},t) \int_{\Sigma} \int_{-1}^{1} G_{\lambda}(x_{\Sigma} + \varepsilon t\nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma}))\nu(y_{\Sigma},s)\Xi(y_{\Sigma},s) \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}).$$

As in the analysis of convergence of $A_{\varepsilon}(\lambda)$, we separately prove the estimates for $||B_{\varepsilon}(\lambda) - \hat{B}_{\varepsilon}(\lambda)||$ and $||\hat{B}_{\varepsilon}(\lambda) - B_{0}(\lambda)||$, which yield then the claimed convergence result. To estimate $||B_{\varepsilon}(\lambda) - \hat{B}_{\varepsilon}(\lambda)||$, we introduce the functions $\tilde{F}_{1,\varepsilon}^{B}, \hat{F}_{1,\varepsilon}^{B}: \Sigma \times (-1, 1) \to [0, \infty]$ by

$$\begin{split} \widetilde{F}_{1,\varepsilon}^{B}(x_{\Sigma},t) &:= \int_{\Sigma} \int_{-1}^{1} \left| u(x_{\Sigma},t) G_{\lambda}(x_{\Sigma} + \varepsilon t \nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) v(y_{\Sigma},s) \right| D_{\varepsilon}(y_{\Sigma},s) \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}), \\ \widehat{F}_{1,\varepsilon}^{B}(y_{\Sigma},s) &:= \int_{\Sigma} \int_{-1}^{1} \left| u(x_{\Sigma},t) G_{\lambda}(x_{\Sigma} + \varepsilon t \nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) v(y_{\Sigma},s) \right| D_{\varepsilon}(y_{\Sigma},s) \mathrm{d}t \mathrm{d}\sigma(x_{\Sigma}). \end{split}$$

Using (3.10) and Proposition A.5 (ii), we get

$$\sup \widetilde{F}_{1,\varepsilon}^B \le c_{B,1}\varepsilon \qquad \text{and} \qquad \sup \widehat{F}_{1,\varepsilon}^B \le c_{B,1}\varepsilon$$

with a constant $c_{B,1} = c_{B,1}(d, \lambda, \Sigma, V) > 0$. Employing these bounds and again applying the Schur test we obtain that

$$\left\|B_{\varepsilon}(\lambda) - \widehat{B}_{\varepsilon}(\lambda)\right\|^{2} \le \left(\sup \widetilde{F}_{1,\varepsilon}^{B}\right) \cdot \left(\sup \widehat{F}_{1,\varepsilon}^{B}\right) \le c_{B,1}^{2}\varepsilon^{2}.$$
(3.13)

To estimate $\|\widehat{B}_{\varepsilon}(\lambda) - B_0(\lambda)\|$ we introduce the auxiliary function $F_{2,\varepsilon}^B \colon \Sigma \times (-1,1) \to [0,\infty]$ as

$$F_{2,\varepsilon}^B(x_{\Sigma},t) := \int_{\Sigma} \int_{-1}^1 \left| u(x_{\Sigma},t) \left(G_{\lambda}(x_{\Sigma} + \varepsilon t \nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) - G_{\lambda}(x_{\Sigma} - y_{\Sigma}) \right) v(y_{\Sigma},y) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}).$$

Using that the absolute value of the integral kernel of $\hat{B}_{\varepsilon}(\lambda) - B_0(\lambda)$ is symmetric and applying Proposition A.9 (ii), we obtain with the help of the Schur test that

$$\left\|\widehat{B}_{\varepsilon}(\lambda) - B_{0}(\lambda)\right\| \leq \sup F_{2,\varepsilon}^{B} \leq c_{B,2}\varepsilon \left(1 + |\ln\varepsilon|\right)$$
(3.14)

with a constant $c_{B,2} = c_{B,2}(d, \lambda, \Sigma, V) > 0$. Putting together the estimates in (3.13), (3.14) and employing the triangle inequality, we eventually deduce that there is a constant $c_B = c_B(d, \lambda, \Sigma, V) > 0$ such that

$$\begin{aligned} \|B_{\varepsilon}(\lambda) - B_{0}(\lambda)\| &\leq \|B_{\varepsilon}(\lambda) - \widehat{B}_{\varepsilon}(\lambda)\| + \|\widehat{B}_{\varepsilon}(\lambda) - B_{0}(\lambda)\| \\ &\leq c_{B,1}\varepsilon + c_{B,2}\varepsilon (1 + |\ln \varepsilon|) \\ &\leq c_{B}\varepsilon (1 + |\ln \varepsilon|). \end{aligned}$$

Finally, we analyze the convergence of $C_{\varepsilon}(\lambda)$. Using again the Schur test, Proposition A.9 (i) and Proposition A.6 we find

$$\begin{aligned} \left\| C_{\varepsilon}(\lambda) - C_{0}(\lambda) \right\|^{2} &\leq \sup_{y \in \mathbb{R}^{d}} \int_{\Sigma} \int_{-1}^{1} \left| u(x_{\Sigma}, t) \left(G_{\lambda}(x_{\Sigma} + \varepsilon t\nu(x_{\Sigma}) - y) - G_{\lambda}(x_{\Sigma} - y) \right) \right| \mathrm{d}t \mathrm{d}\sigma(x_{\Sigma}) \\ &\cdot \sup_{(x_{\Sigma}, t) \in \Sigma \times (-1, 1)} \int_{\mathbb{R}^{d}} \left| u(x_{\Sigma}, t) \left(G_{\lambda}(x_{\Sigma} + \varepsilon t\nu(x_{\Sigma}) - y) - G_{\lambda}(x_{\Sigma} - y) \right) \right| \mathrm{d}t \mathrm{d}\sigma(x_{\Sigma}) \\ &\leq c_{C} \varepsilon^{2} \left(1 + |\ln \varepsilon| \right) \end{aligned}$$

with a constant $c_C = c_C(d, \lambda, \Sigma, V) > 0$. Setting $c := \max\{c_A, c_B, c_C\}$, the claimed result of this lemma follows.

Theorem 2.8, Proposition 3.3 and Lemma 3.4 contain the essential ingredients to prove Theorem 1.1. In the first two steps of the proof the special case Q = 0 is discussed, the general situation is treated in the last step.

Proof of Theorem 1.1. For $\lambda < 0$ and $\varepsilon \in [0, \beta]$ the operators $A_{\varepsilon}(\lambda)$, $B_{\varepsilon}(\lambda)$ and $C_{\varepsilon}(\lambda)$ are defined as in (3.5).

Step 1: First, we prove that $(1 - B_0(\lambda))^{-1}$ exists and is bounded and everywhere defined for $\lambda < 0$ with $|\lambda|$ sufficiently large. Let $M \in (0, 1)$ be fixed and choose $\lambda_M < 0$ such that $||B_{\varepsilon}(\lambda)|| \le M$ for any $\lambda < \lambda_M$ and all $\varepsilon \in (0, \beta]$; recall that such a λ_M exists by Lemma 3.2 (ii). Hence, the operators $(1 - B_{\varepsilon}(\lambda))^{-1}$ are bounded and everywhere defined for $\lambda < \lambda_M$ and $\varepsilon \in (0, \beta]$, and it holds

$$\|(1 - B_{\varepsilon}(\lambda))^{-1}\| \le \frac{1}{1 - M}.$$
(3.15)

Because of (3.15) and Lemma 3.4, we can apply [25, Theorem IV 1.16], which yields that $1 - B_0(\lambda)$ is boundedly invertible. Moreover, for $\lambda < \lambda_M$ we conclude from [25, Theorem IV 1.16] and (3.15) that

$$\left\| (1 - B_{\varepsilon}(\lambda))^{-1} - (1 - B_{0}(\lambda))^{-1} \right\| \leq \frac{\| (1 - B_{\varepsilon}(\lambda))^{-1} \|^{2}}{1 - \| B_{\varepsilon}(\lambda) - B_{0}(\lambda) \| \| (1 - B_{\varepsilon}(\lambda))^{-1} \|} \| B_{\varepsilon}(\lambda) - B_{0}(\lambda) \| \leq c_{1} \varepsilon (1 + |\ln \varepsilon|)$$

holds for all $\varepsilon > 0$ sufficiently small and a constant $c_1 = c_1(d, \lambda, \Sigma, V) > 0$.

Step 2: From Proposition 3.3, Lemma 3.4 and the estimates in Step 1 we obtain

$$\begin{aligned} \left\| (H_{\varepsilon} - \lambda)^{-1} - \left(R(\lambda) + A_0(\lambda)(1 - B_0(\lambda))^{-1}C_0(\lambda) \right) \right\| \\ &= \left\| A_{\varepsilon}(\lambda)(1 - B_{\varepsilon}(\lambda))^{-1}C_{\varepsilon}(\lambda) - A_0(\lambda)(1 - B_0(\lambda))^{-1}C_0(\lambda) \right\| \le c_2 \varepsilon \left(1 + |\ln \varepsilon| \right) \end{aligned}$$

with a constant $c_2 = c_2(d, \lambda, \Sigma, V) > 0$. It remains to verify that

$$R(\lambda) + A_0(\lambda)(1 - B_0(\lambda))^{-1}C_0(\lambda) = (A_{\delta,\alpha} - \lambda)^{-1},$$
(3.16)

where $\alpha \in L^{\infty}(\Sigma)$ is defined as in the theorem,

$$\alpha(x_{\Sigma}) = \int_{-\beta}^{\beta} V(x_{\Sigma} + s\nu(x_{\Sigma})) \mathrm{d}s.$$

In order to show (3.16) let u and v be given by (3.4), introduce the bounded operators $\widehat{U}: L^2(\Sigma) \to L^2(\Sigma \times I^2)$ (-1,1)) and \widehat{V} : $L^2(\Sigma \times (-1,1)) \to L^2(\Sigma)$ by

$$(\widehat{U}\xi)(x_{\Sigma},t) := u(x_{\Sigma},t)\xi(x_{\Sigma}) \text{ and } (\widehat{V}\Xi)(x_{\Sigma}) := \int_{-1}^{1} v(x_{\Sigma},s)\Xi(x_{\Sigma},s)\mathrm{d}s$$

defined almost everywhere, and note that $\widehat{V}\widehat{U}$ is the multiplication operator with α in $L^2(\Sigma)$. Furthermore, recall the definition of the bounded operators $\gamma(\lambda)$, $M(\lambda)$ and the formula for $\gamma(\lambda)^*$ from (2.12); cf. Theorem 2.8. Then, we observe that

$$(A_0(\lambda)\Xi)(x) = \int_{\Sigma} \int_{-1}^{1} G_{\lambda}(x - y_{\Sigma})v(y_{\Sigma}, s)\Xi(y_{\Sigma}, s)\mathrm{d}s\mathrm{d}\sigma(y_{\Sigma})$$
$$= \int_{\Sigma} G_{\lambda}(x - y_{\Sigma}) \bigg(\int_{-1}^{1} v(y_{\Sigma}, s)\Xi(y_{\Sigma}, s)\mathrm{d}s \bigg)\mathrm{d}\sigma(y_{\Sigma})$$
$$= \big(\gamma(\lambda)\widehat{V}\Xi\big)(x)$$

for any $\Xi \in L^2(\Sigma \times (-1, 1))$ and a.e. $x \in \mathbb{R}^d$. Thus, we conclude $A_0(\lambda) = \gamma(\lambda)\widehat{V}$. In a similar way, one finds that

$$(B_0(\lambda)\Xi)(x_{\Sigma},t) = u(x_{\Sigma},t) \int_{\Sigma} \int_{-1}^{1} G_{\lambda}(x_{\Sigma} - y_{\Sigma})v(y_{\Sigma},s)\Xi(y_{\Sigma},s)\mathrm{d}s\mathrm{d}\sigma(y_{\Sigma})$$
$$= u(x_{\Sigma},t) \int_{\Sigma} G_{\lambda}(x_{\Sigma} - y_{\Sigma}) \left(\int_{-1}^{1} v(y_{\Sigma},s)\Xi(y_{\Sigma},s)\mathrm{d}s\right)\mathrm{d}\sigma(y_{\Sigma})$$
$$= \left(\widehat{U}M(\lambda)\widehat{V}\Xi\right)(x_{\Sigma},t)$$

for any $\Xi \in L^2(\Sigma \times (-1,1))$ and a.e. $(x_{\Sigma},t) \in \Sigma \times (-1,1)$, which implies $B_0(\lambda) = \widehat{U}M(\lambda)\widehat{V}$. Finally, one sees that

$$(C_0(\lambda)f)(x_{\Sigma},t) = u(x_{\Sigma},t) \int_{\mathbb{R}^d} G_\lambda(x_{\Sigma}-y)f(y) \mathrm{d}y = \left(\widehat{U}\gamma(\lambda)^*\right)(x_{\Sigma},t)$$

for all $f \in L^2(\mathbb{R}^d)$ and a.e. $(x_{\Sigma}, t) \in \Sigma \times (-1, 1)$, which implies $C_0(\lambda) = \widehat{U}\gamma(\lambda)^*$. Thus, we get

$$\lim_{\varepsilon \to 0} (H_{\varepsilon} - \lambda)^{-1} = R(\lambda) + A_0(\lambda) (1 - B_0(\lambda))^{-1} C_0(\lambda)$$

= $R(\lambda) + \gamma(\lambda) \widehat{V} (1 - \widehat{U}M(\lambda)\widehat{V})^{-1} \widehat{U}\gamma(\lambda)^*.$ (3.17)

Since $\alpha = \widehat{V}\widehat{U}$ we conclude from

$$\widehat{V}\left(1-\widehat{U}M(\lambda)\widehat{V}\right)^{-1}\widehat{U} = \left(1-\widehat{V}\widehat{U}M(\lambda)\right)^{-1}\widehat{V}\widehat{U} = \left(1-\alpha M(\lambda)\right)^{-1}\alpha$$

and (3.17) together with Theorem 2.8 that

$$\lim_{\varepsilon \to 0+} (H_{\varepsilon} - \lambda)^{-1} = R(\lambda) + \gamma(\lambda) (1 - \alpha M(\lambda))^{-1} \alpha \gamma(\lambda)^* = (A_{\delta, \alpha} - \lambda)^{-1}.$$

This completes the proof of Theorem 1.1 in the case Q = 0.

Step 3: Let $Q \in L^{\infty}(\mathbb{R}^d)$ be real valued and let $\lambda \in \mathbb{R}$ be such that $\lambda < \lambda_M - \|Q\|_{L^{\infty}}$. Then λ is smaller than the lower bound of the operators $A_{\delta,\alpha} + Q$ and $H_{\varepsilon} + Q$ for all $\varepsilon \in (0, \beta]$. Using the formula

$$(H_{\varepsilon} + Q - \lambda)^{-1} = \left(1 - (H_{\varepsilon} + Q - \lambda)^{-1}Q\right)(H_{\varepsilon} - \lambda)^{-1},$$

we compute

$$(H_{\varepsilon} + Q - \lambda)^{-1} - (A_{\delta,\alpha} + Q - \lambda)^{-1}$$

$$= \left[(H_{\varepsilon} + Q - \lambda)^{-1} (A_{\delta,\alpha} + Q - \lambda) - 1 \right] (A_{\delta,\alpha} + Q - \lambda)^{-1}$$

$$= \left[(H_{\varepsilon} + Q - \lambda)^{-1} - \left(1 - (H_{\varepsilon} + Q - \lambda)^{-1} Q \right) (A_{\delta,\alpha} - \lambda)^{-1} \right] (A_{\delta,\alpha} - \lambda) (A_{\delta,\alpha} + Q - \lambda)^{-1}$$

$$= \left(1 - (H_{\varepsilon} + Q - \lambda)^{-1} Q \right) \left[(H_{\varepsilon} - \lambda)^{-1} - (A_{\delta,\alpha} - \lambda)^{-1} \right] \left(1 - Q (A_{\delta,\alpha} + Q - \lambda)^{-1} \right).$$

This implies

$$\| (H_{\varepsilon} + Q - \lambda)^{-1} - (A_{\delta,\alpha} + Q - \lambda)^{-1} \|$$

$$\leq \| 1 - (H_{\varepsilon} + Q - \lambda)^{-1} Q \| \cdot \| (H_{\varepsilon} - \lambda)^{-1} - (A_{\delta,\alpha} - \lambda)^{-1} \| \cdot \| 1 - Q (A_{\delta,\alpha} + Q - \lambda)^{-1} \|.$$

Since the norm $\|1 - (H_{\varepsilon} + Q - \lambda)^{-1}Q\|$ is uniformly bounded in ε , the result of Step 2 yields the desired claim.

3.3 Consequences for the spectra of H_{ε} and $A_{\delta,\alpha}$

In this section we discuss how Theorem 1.1 can be used to deduce certain spectral properties of H_{ε} from those of $A_{\delta,\alpha}$ and vice versa.

First, the approximation result in Theorem 1.1 combined with the results in [7, 16, 18] on the existence of geometrically induced bound states for Schrödinger operators with δ -interactions supported on curves and surfaces can be used to show the existence of such bound states also for the operators with regular potentials (in the approximating sequence) provided that the potential well is sufficiently "narrow" and "deep". This application is motivated by Open Problem 7.1 in the review paper [15]. In order to formulate the respective claims we introduce several geometric notions. We say that the hypersurface $\Sigma \subset \mathbb{R}^d$, $d \ge 2$, is a *local deformation* of the hypersurface $\widetilde{\Sigma} \subset \mathbb{R}^d$, if $\Sigma \neq \widetilde{\Sigma}$ and if there exists a compact set $K \subset \mathbb{R}^d$ such that $\Sigma \setminus K = \widetilde{\Sigma} \setminus K$. Furthermore, we introduce for $\theta \in (0, \pi/2)$ the broken line $\mathcal{L} \subset \mathbb{R}^2$ and the infinite circular conical surface $\mathcal{C} \subset \mathbb{R}^3$ by

$$\mathcal{L} := \left\{ (x, y) \in \mathbb{R}^2 \colon y = \cot(\theta) |x| \right\}, \qquad \mathcal{C} := \left\{ (x, y, z) \in \mathbb{R}^3 \colon z = \cot(\theta) \sqrt{x^2 + y^2} \right\}$$

Proposition 3.5 Let $\Sigma \subset \mathbb{R}^d$ be as in Definition 2.1 such that Hypothesis 2.3 holds and assume that Σ satisfies one of the following assumptions.

- (a) d = 2 and Σ is a local deformation of a straight line;
- (b) d = 2 and Σ is a local deformation of a broken line;
- (c) d = 3 and Σ is a C^4 -smooth local deformation of a plane, which admits a global natural parametrization in the sense of [13, Section 2-3, Definition 2];
- (d) d = 3 and Σ is a local deformation of an infinite circular conical surface.

Let the layer $\Omega_{\beta} \subset \mathbb{R}^d$ be as in (2.6) and let $V \in L^{\infty}(\mathbb{R}^d)$ be real valued with $\operatorname{supp} V \subset \Omega_{\beta}$. Assume that $\alpha \in L^{\infty}(\Sigma)$ associated to V as in Theorem 1.1 is a positive constant, which in case (c) is assumed to be sufficiently large. Then for all sufficiently small $\varepsilon > 0$ the self-adjoint operator H_{ε} in (3.1) has a non-empty discrete spectrum below the threshold of its essential spectrum.

Proof. As there is no danger of confusion, we denote the value of the constant positive function $\alpha \in L^{\infty}(\Sigma)$ by α as well. First, recall that by Theorem 1.1 the self-adjoint operators H_{ε} converge in the norm resolvent sense to the self-adjoint lower-semibounded operator $A_{\delta,\alpha}$ as $\varepsilon \to 0+$.

Next, in all the cases it is known that $\sigma_{\text{ess}}(A_{\delta,\alpha}) = [-\alpha^2/4, \infty)$; cf. [6, Theorems 4.2 and 4.10] and [7, Theorem 3.3]. Moreover, Proposition 3.1 implies that the operators H_{ε} are bounded from below uniformly in $\varepsilon \in (0, \beta]$. Hence, the result [38, Satz 9.24 a)] yields

$$\inf \sigma_{\rm ess}(H_{\varepsilon}) \to -\alpha^2/4, \qquad \varepsilon \to 0+.$$
(3.18)

Moreover, in all the cases it is known that $\sigma_d(A_{\delta,\alpha}) \cap (-\infty, -\alpha^2/4) \neq \emptyset$; cf. [16, Theorem 5.2] for (a), (b), [18, Theorem 4.3] for (c) and [7, Theorem 3.3] for (d). Choose now a finite interval $(a, b) \subset (-\infty, -\alpha^2/4)$ with $a, b \in \rho(A_{\delta,\alpha})$ such that $\sigma_d(A_{\delta,\alpha}) \cap (a, b) \neq \emptyset$. By (3.18) we get for all sufficiently small $\varepsilon > 0$ that inf $\sigma_{ess}(H_{\varepsilon}) > b$.

Eventually, it follows from [38, Satz 2.58 a) and Satz 9.24 b)] that for all sufficiently small $\varepsilon > 0$ the dimensions of the spectral subspaces of H_{ε} corresponding to (a, b) are equal to the dimension of the spectral subspace of $A_{\delta,\alpha}$ corresponding to the same interval; i.e.

$$\dim \operatorname{ran} E_{H_{\varepsilon}}((a,b)) = \dim \operatorname{ran} E_{A_{\delta,\alpha}}((a,b)) > 0.$$
(3.19)

Since $b < \inf \sigma_{ess}(H_{\varepsilon})$, this implies the claimed result.

In the next proposition we show that Schrödinger operators with potential wells supported in curved periodic strips have gaps in their spectra under reasonable assumptions. This proposition can be proven in a similar way as Proposition 3.5. It suffices to combine a result [17, Corollary 2.2] on the existence of gaps in the negative spectrum for the Schrödinger operator with a strong δ -interaction supported on a periodic curve with the main result of this paper and with standard statements on spectral convergence. Note that a similar idea was earlier used in a different, albeit related context in [39].

Proposition 3.6 Let $\Sigma \subset \mathbb{R}^2$ be as in Definition 2.1 such that Hypothesis 2.3 holds. Suppose that Σ is parametrized via the tuple $\{\varphi, \mathbb{R}, \mathbb{R}^2\}$ with $|\varphi'| \equiv 1$. For $\varphi(s) = (\varphi_1(s), \varphi_2(s))$ define the signed curvature of Σ by $\kappa := \varphi_1'' \varphi_2' - \varphi_1' \varphi_2''$. Assume that Σ is not a straight line and that κ satisfies the following conditions.

- (a) $\kappa \in C^2(\mathbb{R});$
- (b) $\kappa(s+L) = \kappa(s)$ for some L > 0 and all $s \in \mathbb{R}$;
- (c) $\int_0^L \kappa(s) \mathrm{d}s = 0;$
- (d) $\left|\int_0^T \kappa(s) \mathrm{d}s\right| < \pi/2$ for all $T \in [0, L]$.

Let $\Omega_{\beta} \subset \mathbb{R}^2$ be as in (2.6) and let $V \in L^{\infty}(\mathbb{R}^d)$ be real valued with supp $V \subset \Omega_{\beta}$. Assume that $\alpha \in L^{\infty}(\Sigma)$ associated to V as in Theorem 1.1 is a sufficiently large positive constant. Then for all sufficiently small $\varepsilon > 0$ the self-adjoint operator H_{ε} in (3.1) has at least one spectral gap in the interval $(-\infty, 0)$.

One can also use Theorem 1.1 to obtain spectral results for $A_{\delta,\alpha}$ from those of H_{ε} . To illustrate this idea, we show in the following example that for $d \ge 3$ the operator $A_{\delta,\alpha}$ does not have any negative bound states if Σ is a sphere with radius R > 0 and if $\alpha \in (0, \frac{d-2}{R})$ is a constant. The same result is also contained in [5, Theorem 2.3] for d = 3 and in [4, Theorem 4.1] for arbitrary $d \ge 3$. The proofs there are of a different nature than ours.

Example 3.7 Let $d \ge 3$, R > 0 and $\Sigma = \partial B(0, R)$. Let $\alpha \in (0, \frac{d-2}{R})$ be fixed. Let $q \in C_0^{\infty}(\mathbb{R}_+)$ be non-negative with supp $q \subset (R/2, 3R/2)$ such that

$$\int_{R/2}^{3R/2} q(r) \mathrm{d}r = \alpha$$

Define the radially symmetric potential $V \in L^{\infty}(\mathbb{R}^d)$ by $V(x) := q(|x|), x \in \mathbb{R}^d$. Furthermore, define the scaled potentials V_{ε} and the operators H_{ε} as in the Introduction with $\beta = R/2$ and $\varepsilon \in (0, R/2]$. By Theorem 1.1 the operators H_{ε} converge in the norm resolvent sense to $A_{\delta,\alpha}$ as $\varepsilon \to 0+$. For any $\varepsilon \in (0, R/2]$ the potential V_{ε} is also radially symmetric and we get with the help of Lebesgue's dominated convergence theorem that for $q_{\varepsilon}(|x|) := V_{\varepsilon}(x)$

$$\lim_{\varepsilon \to 0+} \int_0^\infty r q_\varepsilon(r) \mathrm{d}r = \lim_{\varepsilon \to 0+} \int_{-\varepsilon}^\varepsilon (R+r) \frac{R}{2\varepsilon} q\left(R + \frac{Rr}{2\varepsilon}\right) \mathrm{d}r$$
$$= \lim_{\varepsilon \to 0+} \int_{R/2}^{3R/2} \left(\frac{2\varepsilon z}{R} + R - 2\varepsilon\right) q(z) \mathrm{d}z$$
$$= \alpha R < d-2.$$

Hence, using Bargmann's estimate [33, Theorem 3.2], we obtain that H_{ε} has no negative eigenvalues for all sufficiently small $\varepsilon > 0$. Because of Theorem 1.1 and [38, Satz 9.24 a)] it follows that $A_{\delta,\alpha}$ has no negative eigenvalues as well.

A Estimates related to Green's function

In this appendix we provide estimates for integrals that contain Green's function

$$G_{\lambda}(x) = \frac{1}{(2\pi)^{d/2}} \left(\frac{|x|}{\sqrt{-\lambda}}\right)^{1-d/2} K_{d/2-1}\left(\sqrt{-\lambda}|x|\right), \quad x \in \mathbb{R}^d \setminus \{0\}, \quad d \ge 2,$$
(A.1)

from (2.11). The estimates are formulated in a way such that they can be applied directly in the main part of the paper. We note that some of the estimates below are known, but exact references are difficult to find in the

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mathematical literature. Therefore, in order to keep this paper self-contained we also provide complete proofs of standard estimates as, e.g., in Proposition A.1. Throughout this appendix we assume that Σ is an orientable C^2 -hypersurface which satisfies Hypothesis 2.3, and we denote by ν the continuous unit normal vector field of Σ .

In the first preliminary proposition we discuss the asymptotics of G_{λ} and ∇G_{λ} .

Proposition A.1 Let $\lambda \in (-\infty, 0)$ and let G_{λ} be as in (A.1). Then there exists a constant R = R(d) > 0 such that the following assertions hold.

(i) There is a constant c = c(d) such that for all $|x| \leq \frac{R}{\sqrt{-\lambda}}$

$$\left|G_{\lambda}(x)\right| \leq \begin{cases} c\left(1 + \left|\ln\left(\sqrt{-\lambda}|x|\right)\right|\right), & \text{if } d = 2, \\ c|x|^{2-d}, & \text{if } d \geq 3, \end{cases}$$

holds. Moreover, there exists a constant C = C(d) such that for all $|x| > \frac{R}{\sqrt{-\lambda}}$ it holds

$$|G_{\lambda}(x)| \le C \left(\frac{|x|}{\sqrt{-\lambda}}\right)^{1-d/2} e^{-\sqrt{-\lambda}|x|}$$

In particular, $G_{\lambda} \in L^1(\mathbb{R}^d)$.

(ii) The function $\mathbb{R}^d \setminus \{0\} \ni x \mapsto G_\lambda(x)$ is continuously differentiable. Furthermore, there exist constants $c = c(d, \lambda)$ and $C = C(d, \lambda)$ such that for all $|x| \le \frac{R}{\sqrt{-\lambda}}$

$$\left|\nabla G_{\lambda}(x)\right| \le c|x|^{1-d}$$

is true and for all $|x| > \frac{R}{\sqrt{-\lambda}}$ it holds

$$\left|\nabla G_{\lambda}(x)\right| \le Ce^{-\sqrt{-\lambda}|x|}$$

In particular, $\nabla G_{\lambda} \in L^1(\mathbb{R}^d; \mathbb{C}^d)$.

Proof. (i) Due to the asymptotic behavior of $K_{d/2-1}$, see [1, Section 9.6 and 9.7], there exist R = R(d) > 0and constants $C_1, C_2, C_3 > 0$ such that for any $p \in (0, R]$

$$|K_{d/2-1}(p)| \le \begin{cases} C_1(1+|\ln p|), & \text{if } d=2, \\ C_2 p^{1-d/2}, & \text{if } d \ge 3, \end{cases}$$
(A.2)

is satisfied and for any p > R

$$|K_{d/2-1}(p)| \le C_3 e^{-p} \tag{A.3}$$

holds. Hence, the claimed asymptotics follow from the definition of G_{λ} . It is not difficult to check that these asymptotics imply $G_{\lambda} \in L^1(\mathbb{R}^d)$.

(ii) Recall first that the mapping $\mathbb{C} \setminus (-\infty, 0] \ni z \mapsto K_{\nu}(z)$ is holomorphic by [1, Section 9.6]. Therefore, $\mathbb{R}^d \setminus \{0\} \ni x \mapsto G_{\lambda}(x)$ is continuously differentiable and we obtain

$$\nabla G_{\lambda}(x) = \frac{1}{(2\pi)^{d/2}} \frac{x}{|x|} \left(\left(1 - \frac{d}{2} \right) \frac{|x|^{-d/2}}{\sqrt{-\lambda}^{1-d/2}} K_{d/2-1} \left(\sqrt{-\lambda} |x| \right) \right)$$

$$+ \sqrt{-\lambda} \left(\frac{|x|}{\sqrt{-\lambda}} \right)^{1-d/2} K'_{d/2-1} \left(\sqrt{-\lambda} |x| \right) \right)$$
(A.4)

for $d \geq 3$ and

$$\nabla G_{\lambda}(x) = \frac{1}{2\pi} \frac{x}{|x|} \sqrt{-\lambda} K_0'\left(\sqrt{-\lambda}|x|\right)$$
(A.5)

in the case d = 2. Since

$$K'_{d/2-1}(z) = K_{d/2}(z) + \frac{d-2}{2z} K_{d/2-1}(z)$$
(A.6)

by [1, eq. 9.6.26] we conclude from (A.2) that for $d \ge 3$ and $|x| \le \frac{R}{\sqrt{-\lambda}}$

$$|K'_{d/2-1}(\sqrt{-\lambda}|x|)| \le |K_{d/2}(\sqrt{-\lambda}|x|)| + \frac{d-2}{2\sqrt{-\lambda}|x|} |K_{d/2-1}(\sqrt{-\lambda}|x|)| \le C_4|x|^{-d/2} + C_5\frac{d-2}{2\sqrt{-\lambda}|x|}|x|^{1-d/2} \le C_6|x|^{-d/2}$$

holds with some constant $C_6 = C_6(d, \lambda) > 0$. It is easy to see that the same estimate is also true in the case d = 2. Hence, (A.2), (A.4) and (A.5) yield

$$|\nabla G_{\lambda}(x)| \le C_7 \left(|x|^{-d/2} \left| K_{d/2-1} \left(\sqrt{-\lambda} |x| \right) \right| + |x|^{1-d/2} \left| K'_{d/2-1} \left(\sqrt{-\lambda} |x| \right) \right| \right) \le c|x|^{1-d}.$$

In the same way, using (A.3), (A.4), (A.5), and (A.6) one finds for all $|x| \ge \frac{R}{\sqrt{-\lambda}}$ that

$$\left|\nabla G_{\lambda}(x)\right| \leq Ce^{-\sqrt{-\lambda}|x|}$$

holds with some constant $C = C(d, \lambda) > 0$. Finally, it is not difficult to check that the asymptotics of ∇G_{λ} imply $\nabla G_{\lambda} \in L^1(\mathbb{R}^d; \mathbb{C}^d)$.

We start now with the estimates which are needed to show Lemma 3.4. The first proposition in this context is only based on the integrability of G_{λ} ; cf. Proposition A.1 (i).

Proposition A.2 Let Σ be a C^2 -hypersurface which satisfies Hypothesis 2.3 and let $\lambda \in (-\infty, 0)$. Then the following statements are true.

(i) There exists a constant $C = C(d, \lambda) > 0$ such that

$$\sup_{\Sigma \in \Sigma} \int_{\mathbb{R}^d} |G_{\lambda}(x - y_{\Sigma})| \, \mathrm{d}x \le C$$

(ii) Let $\psi \in L^{\infty}(\Sigma \times (-1, 1))$ and let $\varepsilon \in (0, \beta]$. Then there exists a constant $C = C(d, \lambda, \psi) > 0$ such that

$$\sup_{(y_{\Sigma},s)\in\Sigma\times(-1,1)}\int_{\mathbb{R}^d} \left| G_{\lambda}(x-y_{\Sigma}-\varepsilon s\nu(y_{\Sigma}))\psi(y_{\Sigma},s) \right| \mathrm{d}x \leq C.$$

Proof. We only prove item (ii), assertion (i) can be shown in the same way. For $(y_{\Sigma}, s) \in \Sigma \times (-1, 1)$ fixed we have

$$\int_{\mathbb{R}^d} |G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma}))\psi(y_{\Sigma}, s)| \mathrm{d}x \le \|\psi\|_{L^{\infty}} \int_{\mathbb{R}^d} |G_{\lambda}(x)| \mathrm{d}x \le C$$

with some constant $C = C(d, \lambda, \psi)$. This is the claimed result.

The following lemma contains an auxiliary estimate associated to the hypersurface Σ . Recall that σ denotes the Hausdorff measure and Λ_d is the *d*-dimensional Lebesgue measure.

Lemma A.3 Let Σ be a C^2 -hypersurface which satisfies Hypothesis 2.3. Then the following assertions are true.

(i) There exists a constant $\widetilde{C}_1 = \widetilde{C}_1(d, \Sigma) > 0$ such that for any $r_0 > 0$ and any $x \in \mathbb{R}^d$ it holds

$$\sigma(\Sigma \cap B(x, r_0)) \le \widetilde{C}_1 r_0^{d-1}.$$

(ii) Let $\varepsilon \in (0, \beta]$ and let Ω_{ε} be as in (2.6). Then there exists a constant $\widetilde{C}_2 = \widetilde{C}_2(d, \Sigma) > 0$ such that for any $r_0 > 0$ and any $x \in \mathbb{R}^d$ it holds

$$\Lambda_d(\Omega_{\varepsilon} \cap B(x, r_0)) \le \widetilde{C}_2 \varepsilon r_0^{d-1}.$$

Proof. (i) By the definition of the measure σ we have

$$\sigma(\Sigma \cap B(x, r_0)) = \int_{\Sigma \cap B(x, r_0)} \mathrm{d}\sigma(x_{\Sigma}) = \sum_{i \in I} \int_{\varphi_i^{-1}(\Sigma \cap B(x, r_0))} \chi_i(\varphi_i(u)) \sqrt{\det G_i(u)} \mathrm{d}u;$$

cf. (2.2). Assume that $\varphi_i^{-1}(\Sigma \cap B(x, r_0)) \neq \emptyset$, let $u_i \in \varphi_i^{-1}(\Sigma \cap B(x, r_0))$ be fixed and let $u \in \varphi_i^{-1}(\Sigma \cap B(x, r_0))$ be arbitrary. Using Hypothesis 2.3 (c) we obtain

$$c|u-u_i| \le |\varphi_i(u) - \varphi_i(u_i)| \le |\varphi_i(u) - x| + |x - \varphi_i(u_i)| \le 2r_0.$$

Hence, it follows $\varphi_i^{-1}(\Sigma \cap B(x, r_0)) \subset B(u_i, 2r_0/c)$, which implies

$$\sigma(\Sigma \cap B(x, r_0)) = \sum_{i \in I} \int_{\varphi_i^{-1}(\Sigma \cap B(x, r_0))} \chi_i(\varphi_i(u)) \sqrt{\det G_i(u)} du$$
$$\leq C_1 \sum_{i \in I} \int_{\varphi_i^{-1}(\Sigma \cap B(x, r_0))} du$$
$$\leq C_1 \sum_{i \in I} \int_{B(u_i, 2r_0/c)} du \leq \widetilde{C}_1 r_0^{d-1},$$

where we have used that $\det G_i$ is uniformly bounded and that the index set I is finite by assumption.

(ii) Using Proposition 2.6 (ii) and Hypothesis 2.3 (b) we find

$$\begin{split} \Lambda_d(\Omega_{\varepsilon} \cap B(x,r_0)) &= \int_{\Omega_{\varepsilon} \cap B(x,r_0)} \mathrm{d}y \\ &= \int_{\Omega_{\varepsilon}} \mathbb{1}_{\Omega_{\varepsilon} \cap B(x,r_0)}(y) \mathrm{d}y \\ &= \int_{\Sigma} \int_{-\varepsilon}^{\varepsilon} \mathbb{1}_{\Omega_{\varepsilon} \cap B(x,r_0)}(y_{\Sigma} + s\nu(y_{\Sigma})) \det(1 - sW(y_{\Sigma})) \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ &\leq C_2 \int_{\Sigma} \int_{-\varepsilon}^{\varepsilon} \mathbb{1}_{\Omega_{\varepsilon} \cap B(x,r_0)}(y_{\Sigma} + s\nu(y_{\Sigma})) \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}). \end{split}$$

Let ι_{φ_i} be given by (2.3). Using the definition of the measure σ and the fact that det G_i is uniformly bounded by assumption, it follows

$$\begin{split} \Lambda_d(\Omega_{\varepsilon} \cap B(x,r_0)) &\leq C_2 \int_{\Sigma} \int_{-\varepsilon}^{\varepsilon} \mathbbm{1}_{\Omega_{\varepsilon} \cap B(x,r_0)} (y_{\Sigma} + s\nu(y_{\Sigma})) \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ &= C_2 \sum_{i \in I} \int_{U_i} \int_{-\varepsilon}^{\varepsilon} \chi_i(\varphi_i(u)) \mathbbm{1}_{\Omega_{\varepsilon} \cap B(x,r_0)} (\iota_{\varphi_i}(u,s)) \sqrt{\det G_i(u)} \mathrm{d}s \mathrm{d}u \\ &\leq C_3 \sum_{i \in I} \int_{U_i} \int_{-\varepsilon}^{\varepsilon} \mathbbm{1}_{A_i}(u,s) \mathrm{d}s \mathrm{d}u, \end{split}$$

where $A_i := \iota_{\varphi_i}^{-1}(\Omega_{\varepsilon} \cap B(x, r_0)) \subset U_i \times (-\varepsilon, \varepsilon)$ and hence $\mathbb{1}_{A_i} = \mathbb{1}_{\Omega_{\varepsilon} \cap B(x, r_0)} \circ \iota_{\varphi_i}$. Assume $A_i \neq \emptyset$, let $(u_i, t) \in A_i$ be fixed and let $(u, s) \in A_i$ be arbitrary. Using Hypothesis 2.3 (c) we find

$$\begin{aligned} c|u - u_i| &\leq c \left(|u - u_i|^2 + (s - t)^2 \right)^{1/2} \\ &\leq |\varphi_i(u) + s\nu(\varphi_i(u)) - \varphi_i(u_i) - t\nu(\varphi_i(u_i))| \\ &\leq |\varphi_i(u) + s\nu(\varphi_i(u)) - x| + |x - \varphi_i(u_i) - t\nu(\varphi_i(u_i))| \\ &< 2r_0. \end{aligned}$$

Therefore, we obtain $A_i \subset B(u_i, 2r_0/c) \times (-\varepsilon, \varepsilon)$, which yields

$$\Lambda_d(\Omega_{\varepsilon} \cap B(x, r_0)) \le C_3 \sum_{i \in I} \int_{U_i} \int_{-\varepsilon}^{\varepsilon} \mathbb{1}_{A_i}(u, s) ds du$$
$$\le C_3 \sum_{i \in I} \int_{B(u_i, 2r_0/c)} \int_{-\varepsilon}^{\varepsilon} ds du$$
$$= \widetilde{C}_2 \varepsilon r_0^{d-1}.$$

The next proposition contains two estimates of a similar flavour as in Proposition A.2. The proof is essentially based on Lemma A.3.

Proposition A.4 Let Σ be a C^2 -hypersurface which satisfies Hypothesis 2.3 and let $\lambda \in (-\infty, 0)$. Then the following assertions are true.

(i) There is a constant $C = C(d, \lambda, \Sigma) > 0$ such that

$$\sup_{x \in \mathbb{R}^d} \int_{\Sigma} \left| G_{\lambda}(x - y_{\Sigma}) \right| \mathrm{d}\sigma(y_{\Sigma}) \leq C.$$

Moreover, for any fixed C > 0 there exists $\lambda_C < 0$ such that the above inequality is satisfied for this fixed C and all $\lambda < \lambda_C$.

(ii) Let $\varepsilon \in (0, \beta]$ and let Ω_{ε} be as in (2.6). Then there is a constant $M = M(d, \lambda, \Sigma) > 0$ such that

$$\sup_{x \in \mathbb{R}^d} \int_{\Omega_{\varepsilon}} |G_{\lambda}(x-y)| \mathrm{d} y \le M \varepsilon.$$

Moreover, for any fixed M > 0 there exists $\lambda_M < 0$ such that the above inequality is satisfied for this fixed M, all $\lambda < \lambda_M$ and all $\varepsilon \in (0, \beta]$.

Proof. (i) Let $x \in \mathbb{R}^d$ and $\lambda \in (-\infty, 0)$ be fixed, let R > 0 be as in Proposition A.1 and set

$$\Sigma_1 = \{y_{\Sigma} \in \Sigma : \sqrt{-\lambda}|x - y_{\Sigma}| < R\}$$
 and $\Sigma_2 = \{y_{\Sigma} \in \Sigma : \sqrt{-\lambda}|x - y_{\Sigma}| > R\}.$

Then we have

$$\int_{\Sigma} |G_{\lambda}(x-y_{\Sigma})| \mathrm{d}\sigma(y_{\Sigma}) = \int_{\Sigma_1} |G_{\lambda}(x-y_{\Sigma})| \mathrm{d}\sigma(y_{\Sigma}) + \int_{\Sigma_2} |G_{\lambda}(x-y_{\Sigma})| \mathrm{d}\sigma(y_{\Sigma}).$$

In order to find an estimate for the integral over Σ_1 observe that

$$\int_{\Sigma_1} |G_{\lambda}(x - y_{\Sigma})| \mathrm{d}\sigma(y_{\Sigma}) = \sum_{n=1}^{\infty} \int_{A_n} |G_{\lambda}(x - y_{\Sigma})| \mathrm{d}\sigma(y_{\Sigma}),$$

where the sets A_n are defined as

$$A_n = \left\{ y_{\Sigma} \in \Sigma : 2^{-n} < \sqrt{-\lambda} |x - y_{\Sigma}| / R < 2^{-n+1} \right\}, \qquad n \in \mathbb{N}.$$

Due to the asymptotics of G_λ in Proposition A.1 (i) we find for $y_\Sigma \in A_n$ that

$$|G_{\lambda}(x - y_{\Sigma})| \le \begin{cases} C_1 (1 + n \ln 2), & \text{if } d = 2, \\ C_2 (-\lambda)^{d/2 - 1} \cdot (2^{d-2})^n, & \text{if } d \ge 3. \end{cases}$$

Hence, we get in the case $d \ge 3$ that

$$\int_{\Sigma_1} |G_\lambda(x - y_\Sigma)| \mathrm{d}\sigma(y_\Sigma) = \sum_{n=1}^\infty \int_{A_n} |G_\lambda(x - y_\Sigma)| \mathrm{d}\sigma(y_\Sigma)$$
$$\leq \sum_{n=1}^\infty C_2 (-\lambda)^{d/2 - 1} \cdot (2^{d-2})^n \int_{A_n} \mathrm{d}\sigma(y_\Sigma)$$

Using that $A_n \subset B(x, R \cdot 2^{-n+1}/\sqrt{-\lambda})$, we can employ Lemma A.3 (i) and get

$$\int_{\Sigma_{1}} |G_{\lambda}(x - y_{\Sigma})| d\sigma(y_{\Sigma}) \leq C_{3}(-\lambda)^{d/2 - 1} \sum_{n=1}^{\infty} (2^{d-2})^{n} \cdot (R \cdot 2^{-n+1}/\sqrt{-\lambda})^{d-1} \\ \leq C_{4}(-\lambda)^{-1/2} \sum_{n=1}^{\infty} 2^{-n} \\ = C_{4}(-\lambda)^{-1/2}.$$
(A.7)

Note that similar estimates are also true in the case d = 2, as for $y_{\Sigma} \in A_n$ it holds

$$|G_{\lambda}(x-y_{\Sigma})| \le C_1(1+n\ln 2) \le C_5(1+2^{n/2}) \le 2C_5 2^{n/2}.$$

Finally, we derive an estimate for the integral over Σ_2 . For this purpose we decompose the integral as follows

$$\int_{\Sigma_2} |G_{\lambda}(x-y_{\Sigma})| \mathrm{d}\sigma(y_{\Sigma}) = \sum_{n=1}^{\infty} \int_{B_n} |G_{\lambda}(x-y_{\Sigma})| \mathrm{d}\sigma(y_{\Sigma}),$$

where the sets B_n are given by

$$B_n = \left\{ y_{\Sigma} \in \Sigma : R + n - 1 < \sqrt{-\lambda} |x - y_{\Sigma}| < R + n \right\}, \qquad n \in \mathbb{N}.$$

Using the asymptotics of G_{λ} for large arguments from Proposition A.1 (i) we find for $y_{\Sigma} \in B_n$ that

$$\left|G_{\lambda}(x-y_{\Sigma})\right| \leq C_6 \left(\frac{|x-y_{\Sigma}|}{\sqrt{-\lambda}}\right)^{1-d/2} e^{-\sqrt{-\lambda}|x-y_{\Sigma}|} \leq C_7(-\lambda)^{d/2-1} e^{-(R+n-1)}.$$

Since $B_n \subset B(x, (R+n)/\sqrt{-\lambda})$, we can employ Lemma A.3 (i) and get that

$$\int_{\Sigma_2} |G_\lambda(x-y_\Sigma)| \mathrm{d}\sigma(y_\Sigma) = \sum_{n=1}^{\infty} \int_{B_n} |G_\lambda(x-y_\Sigma)| \mathrm{d}\sigma(y_\Sigma)$$

$$\leq \sum_{n=1}^{\infty} C_7(-\lambda)^{d/2-1} e^{-(R+n-1)} \int_{B_n} \mathrm{d}\sigma(y_\Sigma)$$

$$\leq C_8(-\lambda)^{-1/2} \sum_{n=1}^{\infty} e^{-n} (R+n)^{d-1}$$

$$= C_9(-\lambda)^{-1/2}.$$
(A.8)

The estimates (A.7) and (A.8) yield now the claimed bounds. In particular, by choosing $\lambda < 0$ with $|\lambda|$ sufficiently large the constant $C := (C_4 + C_9)(-\lambda)^{-1/2}$ becomes arbitrarily small.

(ii) The proof of this statement is similar to the one of assertion (i). The difference is to replace Σ by the layer Ω_{ε} and to use Lemma A.3 (ii) instead of Lemma A.3 (i).

The next result is a consequence of Proposition A.4.

Proposition A.5 Let Σ be a C^2 -hypersurface satisfying Hypothesis 2.3, let $\lambda \in (-\infty, 0)$ and let $\varepsilon \in (0, \beta]$. Then the following statements are true.

(i) Let $\psi \in L^{\infty}(\Sigma \times (-1, 1))$. Then there exists a constant $C = C(d, \lambda, \Sigma, \psi) > 0$ such that

$$\sup_{x \in \mathbb{R}^d} \int_{\Sigma} \int_{-1}^{1} \left| G_{\lambda}(x - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) \psi(y_{\Sigma}, s) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \le C.$$

(ii) Let $\omega, \psi \in L^{\infty}(\Sigma \times (-1, 1))$. Then there exists a constant $C = C(d, \lambda, \Sigma, \omega, \psi) > 0$ such that

$$\sup_{(x_{\Sigma},t)\in\Sigma\times(-1,1)}\int_{\Sigma}\int_{-1}^{1}|\omega(x_{\Sigma},t)G_{\lambda}(x_{\Sigma}+\varepsilon t\nu(x_{\Sigma})-y_{\Sigma}-\varepsilon s\nu(y_{\Sigma}))\psi(y_{\Sigma},s)|\mathrm{d}s\mathrm{d}\sigma(y_{\Sigma})\leq C.$$

Proof. Let $z \in \mathbb{R}^d$ be fixed. According to Hypothesis 2.3 (b) there exists a constant $C_1 > 0$ such that for any $s \in (-1, 1)$ and all $y_{\Sigma} \in \Sigma$ the estimate

$$1 \le C_1 \det(1 - \varepsilon s W(y_{\Sigma}))$$

holds. Using Proposition 2.6 (ii) and Proposition A.4 (ii) we conclude that there exists a constant C > 0 (independent of z) such that

$$\begin{split} \int_{\Sigma} \int_{-1}^{1} \left| G_{\lambda}(z - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ &\leq C_{1} \int_{\Sigma} \int_{-1}^{1} \left| G_{\lambda}(z - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) \right| \mathrm{det}(1 - \varepsilon s W(y_{\Sigma})) \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ &= \frac{C_{1}}{\varepsilon} \int_{\Sigma} \int_{-\varepsilon}^{\varepsilon} \left| G_{\lambda}(z - y_{\Sigma} - r \nu(y_{\Sigma})) \right| \mathrm{det}(1 - r W(y_{\Sigma})) \mathrm{d}r \mathrm{d}\sigma(y_{\Sigma}) \\ &= \frac{C_{1}}{\varepsilon} \int_{\Omega_{\varepsilon}} \left| G_{\lambda}(z - y) \right| \mathrm{d}y \\ &\leq C. \end{split}$$
(A.9)

(i) Let $x \in \mathbb{R}^d$ be fixed. Then it follows from (A.9) (with z = x) that

$$\begin{split} \int_{\Sigma} \int_{-1}^{1} \left| G_{\lambda}(x - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) \psi(y_{\Sigma}, s) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ & \leq \|\psi\|_{L^{\infty}} \int_{\Sigma} \int_{-1}^{1} \left| G_{\lambda}(x - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ & \leq C, \end{split}$$

where C does not depend on x. This is the claimed result.

(ii) Let $x_{\Sigma} \in \Sigma$ and $t \in (-1, 1)$ be fixed. Then it follows from (A.9) (with $z = x_{\Sigma} + \varepsilon t \nu(x_{\Sigma})$) that

$$\begin{split} \int_{\Sigma} \int_{-1}^{1} \left| \omega(x_{\Sigma}, t) G_{\lambda}(x_{\Sigma} + \varepsilon t \nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) \psi(y_{\Sigma}, s) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ & \leq \|\omega\|_{L^{\infty}} \|\psi\|_{L^{\infty}} \int_{\Sigma} \int_{-1}^{1} \left| G_{\lambda}(x_{\Sigma} + \varepsilon t \nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ & \leq C, \end{split}$$

where C does not depend on x_{Σ} and t. This completes the proof of Proposition A.5.

The following estimates are slightly more involved than the previous ones, as from now on also the gradient of G_{λ} has to be considered. First, we provide a useful simple argument that is needed for the next results. Assume that $\tau \in L^{\infty}(\mathbb{R}^d \times \Sigma \times (-1, 1)^2; \mathbb{R}^d)$ and that $(x, y_{\Sigma}, s, t) \in \mathbb{R}^d \times \Sigma \times (-1, 1)^2$ is such that

$$x - y_{\Sigma} + \varepsilon \theta \tau(x, y_{\Sigma}, s, t) \neq 0$$

for all $\theta \in [0,1]$ and all $\varepsilon \in (0,\beta]$. Since $G_{\lambda}(y)$ is differentiable for $y \neq 0$ (see Proposition A.1 (ii)) it follows that the mapping $[0,1] \ni \theta \mapsto G_{\lambda}(x - y_{\Sigma} + \varepsilon \theta \tau(x, y_{\Sigma}, s, t))$ is differentiable and one has

$$\begin{aligned} G_{\lambda}(x - y_{\Sigma} + \varepsilon \tau(x, y_{\Sigma}, s, t)) &- G_{\lambda}(x - y_{\Sigma}) \\ &= \int_{0}^{1} \frac{\mathrm{d}}{\mathrm{d}\theta} G_{\lambda}(x - y_{\Sigma} + \varepsilon \theta \tau(x, y_{\Sigma}, s, t)) \mathrm{d}\theta \\ &= \int_{0}^{1} \left\langle \nabla G_{\lambda}(x - y_{\Sigma} + \varepsilon \theta \tau(x, y_{\Sigma}, s, t)), \varepsilon \tau(x, y_{\Sigma}, s, t) \right\rangle \mathrm{d}\theta. \end{aligned}$$

Using the Cauchy-Schwarz inequality for vectors in \mathbb{R}^d this leads to

$$\begin{aligned} \left| G_{\lambda}(x - y_{\Sigma} + \varepsilon \tau(x, y_{\Sigma}, s, t)) - G_{\lambda}(x - y_{\Sigma}) \right| \\ &\leq \varepsilon \|\tau\|_{L^{\infty}} \int_{0}^{1} \left| \nabla G_{\lambda}(x - y_{\Sigma} + \varepsilon \theta \tau(x, y_{\Sigma}, s, t)) \right| \mathrm{d}\theta. \end{aligned}$$
(A.10)

Proposition A.6 Let Σ be a C^2 -hypersurface which satisfies Hypothesis 2.3, let $\psi \in L^{\infty}(\Sigma \times (-1,1))$ and let $\lambda \in (-\infty, 0)$. Then there exists a constant $C = C(d, \lambda, \psi) > 0$ such that for all $\varepsilon \in (0, \beta]$ it holds

$$\sup_{(y_{\Sigma},s)\in\Sigma\times(-1,1)}\int_{\mathbb{R}^d} \left| \left(G_{\lambda}(x-y_{\Sigma}-\varepsilon s\nu(y_{\Sigma})) - G_{\lambda}(x-y_{\Sigma}) \right) \psi(y_{\Sigma},s) \right| \mathrm{d}x \le C\varepsilon$$

Proof. Let $y_{\Sigma} \in \Sigma$ and $s \in (-1, 1)$ be fixed. Since $x - y_{\Sigma} - \varepsilon s \theta \nu(y_{\Sigma}) \neq 0$ for a.e. $x \in \mathbb{R}^d$, it follows from (A.10) (with $\tau(x, y_{\Sigma}, s, t) = -s\nu(y_{\Sigma})$ implying $\|\tau\|_{L^{\infty}} = 1$)

$$\begin{split} \int_{\mathbb{R}^d} \left| \left(G_\lambda(x - y_\Sigma - \varepsilon s\nu(y_\Sigma)) - G_\lambda(x - y_\Sigma) \right) \psi(y_\Sigma, s) \right| \mathrm{d}x \\ &\leq \varepsilon \|\psi\|_{L^\infty} \int_{\mathbb{R}^d} \int_0^1 \left| \nabla G_\lambda(x - y_\Sigma - \varepsilon s\theta\nu(y_\Sigma)) \right| \mathrm{d}\theta \mathrm{d}x. \end{split}$$

Now consider the bijective transformation $T: \mathbb{R}^d \times (0,1) \to \mathbb{R}^d \times (0,1)$ given by

$$\begin{pmatrix} \xi \\ \phi \end{pmatrix} = T(x,\theta) := \begin{pmatrix} x - y_{\Sigma} - \varepsilon s \theta \nu(y_{\Sigma}) \\ \theta \end{pmatrix}$$

Note that T is differentiable and that its Jacobian matrix is given by

$$DT(x,\theta) = \begin{pmatrix} \mathbb{1}_d & -\varepsilon s\nu(y_{\Sigma}) \\ 0 & 1 \end{pmatrix},$$

where $\mathbb{1}_d$ is the identity matrix in $\mathbb{R}^{d \times d}$. Hence $|\det DT(x, \theta)| = 1$ and we conclude that

$$\begin{split} \int_{\mathbb{R}^d} \left| \left(G_\lambda(x - y_\Sigma - \varepsilon s\nu(y_\Sigma)) - G_\lambda(x - y_\Sigma) \right) \psi(y_\Sigma, s) \right| \mathrm{d}x \\ &\leq \varepsilon \|\psi\|_{L^\infty} \int_{\mathbb{R}^d} \int_0^1 \left| \nabla G_\lambda(x - y_\Sigma - \varepsilon s\theta\nu(y_\Sigma)) \right| \mathrm{d}\theta \mathrm{d}x \\ &= \varepsilon \|\psi\|_{L^\infty} \int_{\mathbb{R}^d} \int_0^1 \left| \nabla G_\lambda(\xi) \right| \mathrm{d}\phi \mathrm{d}\xi \\ &= \varepsilon \|\psi\|_{L^\infty} \int_{\mathbb{R}^d} \left| \nabla G_\lambda(\xi) \right| \mathrm{d}\xi, \end{split}$$

where we used in the last step that the integrand was independent of ϕ . Since the last integral is finite by Proposition A.1 (ii), the proof is complete.

The next lemma contains an auxiliary estimate that is needed to prove the final two integral bounds.

Lemma A.7 Let Σ be a C^2 -smooth hypersurface satisfying Hypothesis 2.3, let $\lambda \in (-\infty, 0)$, let $r_0 > 0$ and let $\widetilde{\Sigma} \subset \Sigma$ be measurable. Then there exists a constant $C = C(r_0, d, \lambda, \Sigma) > 0$ such that

$$\int_{\widetilde{\Sigma}} \int_{-1}^{1} \left| \nabla G_{\lambda} (x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \le C$$

for all sufficiently small $\varepsilon > 0$, all $\theta \in (0,1)$ and all $x \in \mathbb{R}^d$ satisfying $\operatorname{dist}(x, \widetilde{\Sigma}) \ge r_0$.

Proof. Let $x \in \mathbb{R}^d$ with $\operatorname{dist}(x, \widetilde{\Sigma}) \ge r_0$ and $\theta \in (0, 1)$ be fixed. We are going to show first that

$$\int_{\widetilde{\Sigma}} \left| \nabla G_{\lambda} (x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}\sigma(y_{\Sigma}) \le \widetilde{C}$$
(A.11)

with \widetilde{C} independent of x, s, θ and ε . For this define the sets

$$A_n := \left\{ y_{\Sigma} \in \widetilde{\Sigma} : r_0 + n - 1 < |x - y_{\Sigma}| < r_0 + n \right\}, \qquad n \in \mathbb{N}.$$

Since it holds for $\varepsilon \leq \frac{r_0}{2}$ and all $(y_{\Sigma}, s) \in \widetilde{\Sigma} \times (-1, 1)$

$$|x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})| \ge |x - y_{\Sigma}| - \varepsilon \ge \frac{r_0}{2},$$

it follows from Proposition A.1 (ii) that for $y_{\Sigma} \in A_n \subset \widetilde{\Sigma}$

$$\begin{aligned} \left| \nabla G_{\lambda} (x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| &\leq C_1 e^{-\sqrt{-\lambda} |x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})|} \\ &\leq C_2 e^{-\sqrt{-\lambda} |x - y_{\Sigma}|} \\ &\leq C_2 e^{-\sqrt{-\lambda} (r_0 + n - 1)}. \end{aligned}$$

Finally, since $A_n \subset B(x, r_0 + n)$ we can employ Lemma A.3 (i) and get that

$$\begin{split} \int_{\widetilde{\Sigma}} \left| \nabla G_{\lambda} \left(x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma}) \right) \right| \mathrm{d}\sigma(y_{\Sigma}) &= \sum_{n=1}^{\infty} \int_{A_n} \left| \nabla G_{\lambda} \left(x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma}) \right) \right| \mathrm{d}\sigma(y_{\Sigma}) \\ &\leq C_2 \sum_{n=1}^{\infty} \int_{A_n} e^{-\sqrt{-\lambda}(r_0 + n - 1)} \mathrm{d}\sigma(y_{\Sigma}) \\ &\leq C_3 \sum_{n=1}^{\infty} e^{-\sqrt{-\lambda} \cdot n} (r_0 + n)^{d-1} \\ &= C_4. \end{split}$$

Thus, the estimate (A.11) is true. Finally, interchanging the order of integration we obtain

$$\int_{\widetilde{\Sigma}} \int_{-1}^{1} \left| \nabla G_{\lambda}(x - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \leq \int_{-1}^{1} C_4 \mathrm{d}s \leq C.$$

This is the claimed result.

The next lemma contains the main tool to prove the final two integral bounds. Here, we apply Lemma A.7 for $\widetilde{\Sigma} = \Sigma$ and $\widetilde{\Sigma} = \Sigma \setminus B(x, r_0)$ with $x \in \mathbb{R}^d$. Moreover, we make use of the notation $C([0, 1]; \mathbb{R}^d)$ for the set of all continuous vector valued functions defined on the interval [0, 1].

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Lemma A.8 Let Σ be a C^2 -smooth hypersurface satisfying Hypothesis 2.3 and let $\lambda \in (-\infty, 0)$. Then there exists a constant $C = C(d, \lambda, \Sigma) > 0$ such that for all $\varepsilon > 0$ sufficiently small and all $x \in C([0, 1]; \mathbb{R}^d)$ it holds

$$\int_{\Sigma} \int_{-1}^{1} \int_{0}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}\theta \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \leq C \left(1 + \left| \ln \varepsilon \right| \right).$$

Proof. Let $x \in C([0,1]; \mathbb{R}^d)$ be fixed and choose $r_0 > 0$. Interchanging the order of integration we get

$$\begin{split} \int_{\Sigma} \int_{-1}^{1} \int_{0}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}\theta \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ &= \int_{0}^{1} \int_{\Sigma} \int_{-1}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \mathrm{d}\theta. \end{split}$$

We are going to find a suitable bound for the integral with respect to ds and $d\sigma$. So let $\theta \in (0, 1)$ be fixed. If $dist(x(\theta), \Sigma) > r_0$, Lemma A.7 yields immediately that

$$\int_{\Sigma} \int_{-1}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \leq C_{1}$$

for all $\varepsilon > 0$ sufficiently small and thus, the claimed result is true in this case.

If $dist(x(\theta), \Sigma) \leq r_0$ we split the hypersurface Σ into two disjoint parts

$$\Sigma_1 := \Sigma \cap B(x(\theta), r_0)$$
 and $\Sigma_2 := \Sigma \setminus B(x(\theta), r_0), \quad \Sigma = \Sigma_1 \cup \Sigma_2,$

and we define the following two auxiliary quantities

$$\mathcal{I}_{j}(x(\theta), \theta, \varepsilon) := \int_{\Sigma_{j}} \int_{-1}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}), \quad j \in \{1, 2\}$$

Hence, it holds that

$$\int_{\Sigma} \int_{-1}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| ds d\sigma(y_{\Sigma}) = \mathcal{I}_{1}(x(\theta), \theta, \varepsilon) + \mathcal{I}_{2}(x(\theta), \theta, \varepsilon).$$
(A.12)

Again Lemma A.7 implies that $\mathcal{I}_2(x(\theta), \theta, \varepsilon) \leq C_1$ for all sufficiently small $\varepsilon > 0$ independent of $x(\theta)$ and θ .

It remains to estimate $\mathcal{I}_1(x(\theta), \theta, \varepsilon)$. Let $\{\varphi_i, U_i, V_i\}_{i \in I}$ be a parametrization of Σ as in Definition 2.1. By (2.2) we get

$$\mathcal{I}_1(x(\theta), \theta, \varepsilon) = \sum_{i \in I} \int_{-1}^1 \int_{\varphi_i^{-1}(\Sigma_1)} \chi_i(\varphi_i(v)) \left| \nabla G_\lambda \left(x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v)) \right) \right| \sqrt{\det G_i(v)} dv ds,$$

where $\{\chi_i\}_{i\in I}$ is a partition of unity subordinate to $\{V_i\}_{i\in I}$. Since det $G_i(v)$ is bounded by Definition 2.1 (e), we can continue estimating

$$\begin{aligned} \mathcal{I}_{1}(x(\theta),\theta,\varepsilon) &= \sum_{i\in I} \int_{-1}^{1} \int_{\varphi_{i}^{-1}(\Sigma_{1})} \chi_{i}(\varphi_{i}(v)) \left| \nabla G_{\lambda} \left(x(\theta) - \varphi_{i}(v) - \varepsilon \theta s \nu(\varphi_{i}(v)) \right) \right| \sqrt{\det G_{i}(v)} \mathrm{d}v \mathrm{d}s \\ &\leq C_{2} \sum_{i\in I} \int_{-1}^{1} \int_{\varphi_{i}^{-1}(\Sigma_{1})} \left| \nabla G_{\lambda} \left(x(\theta) - \varphi_{i}(v) - \varepsilon \theta s \nu(\varphi_{i}(v)) \right) \right| \mathrm{d}v \mathrm{d}s \\ &\leq C_{3} \sum_{i\in I} \int_{-1}^{1} \int_{\varphi_{i}^{-1}(\Sigma_{1})} \left| x(\theta) - \varphi_{i}(v) - \varepsilon \theta s \nu(\varphi_{i}(v)) \right|^{1-d} \mathrm{d}v \mathrm{d}s, \end{aligned}$$

(A.13)

where we used an estimate for $|\nabla G_{\lambda}|$ that follows from Proposition A.1 (ii). It remains to find a suitable bound for $|x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))|^{1-d}$.

Claim. Let $i \in I$ with $\varphi_i^{-1}(\Sigma_1) \neq \emptyset$ be fixed. Then there exists a constant $\widetilde{C} = \widetilde{C}(i, \Sigma)$ and parameters $u = u(x(\theta), \theta, \varepsilon, i) \in \overline{U}_i$ and $t = t(x(\theta), \theta, \varepsilon, i) \in [-1, 1]$ with the following properties:

(a) for any $(v,s) \in \varphi_i^{-1}(\Sigma_1) \times (-1,1)$ it holds

$$|x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))|^{1-d} \le \widetilde{C}(|u-v|^{d-1} + (\varepsilon \theta |t-s|)^{d-1})^{-1};$$

(b) $\varphi_i^{-1}(\Sigma_1) \subset B(u, 2r_0/c).$

Here, c is the same constant as in Hypothesis 2.3 (c). We mention that u and t also depend on d, Σ and r_0 , but this is not of importance in the following.

In order to prove the claim, set

$$\mathsf{h} := \inf \big\{ |x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))| \colon (v, s) \in \varphi_i^{-1}(\Sigma_1) \times (-1, 1) \big\}.$$

Choose a sequence $(u_n, t_n) \in \varphi_i^{-1}(\Sigma_1) \times (-1, 1)$ such that

$$|x(\theta) - \varphi_i(u_n) - \varepsilon \theta t_n \nu(\varphi_i(u_n))| \to \mathsf{h} \quad \text{and} \quad \varphi_i(u_n) + \varepsilon \theta t_n \nu(\varphi_i(u_n)) \to y \in \mathbb{R}^d;$$

this is possible since the set $\{\varphi_i(v) + \varepsilon \theta s \nu(\varphi_i(v)) \colon (v, s) \in \varphi_i^{-1}(\Sigma_1) \times (-1, 1)\}$ is bounded in \mathbb{R}^d . Using Hypothesis 2.3 (c), we find that

$$c^{2}(|u_{n}-u_{m}|^{2}+(\varepsilon\theta)^{2}|t_{n}-t_{m}|^{2}) \leq |\varphi_{i}(u_{n})+\varepsilon\theta t_{n}\nu(\varphi_{i}(u_{n}))-\varphi_{i}(u_{m})-\varepsilon\theta t_{m}\nu(\varphi_{i}(u_{m}))|^{2}.$$

Therefore, (u_n) and (t_n) are Cauchy sequences. Set $\lim u_n =: u \in \overline{U_i}$ and $\lim t_n =: t \in [-1, 1]$. Using a continuity argument and again Hypothesis 2.3 (c), we find for all $v \in U_i$ and all $s \in (-1, 1)$ that

$$c^{2}(|u-v|^{2}+(\varepsilon\theta)^{2}|t-s|^{2}) = c^{2} \lim_{n \to \infty} \left(|u_{n}-v|^{2}+(\varepsilon\theta)^{2}|t_{n}-s|^{2}\right)$$

$$\leq \lim_{n \to \infty} |\varphi_{i}(u_{n})+\varepsilon\theta t_{n}\nu(\varphi_{i}(u_{n}))-\varphi_{i}(v)-\varepsilon\theta s\nu(\varphi_{i}(v))|^{2} \quad (A.14)$$

$$= |y-\varphi_{i}(v)-\varepsilon\theta s\nu(\varphi_{i}(v))|^{2}.$$

Finally, for $v \in \varphi_i^{-1}(\Sigma_1)$ and $s \in (-1, 1)$ it holds by the triangle inequality

$$\begin{aligned} |y - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))| &\leq |x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))| + |x(\theta) - y| \\ &\leq 2|x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))| \end{aligned}$$

due to the construction of y. Hence, we get that

$$|x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))| \ge \frac{1}{2} |y - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v))|.$$
(A.15)

In particular, this and (A.14) with s = 0 imply that

$$c|u-v| \le c(|u-v|^2 + (\varepsilon\theta t)^2)^{1/2} \le |y-\varphi_i(v)| \le 2|x(\theta) - \varphi_i(v)| \le 2r_0$$

and thus, $\varphi_i^{-1}(\Sigma_1) \subset B(u, 2r_0/c)$. This is item (b) of the claim. Furthermore, using for a, b > 0 the inequality

 $a^{d-1} + b^{d-1} \leq (a+b)^{d-1} \quad \Leftrightarrow \quad (a+b)^{1-d} \leq \left(a^{d-1} + b^{d-1}\right)^{-1}$

and equations (A.14) and (A.15) we obtain

$$\begin{aligned} \left| x(\theta) - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v)) \right|^{1-d} &\leq 2^{d-1} \left| y - \varphi_i(v) - \varepsilon \theta s \nu(\varphi_i(v)) \right|^{1-d} \\ &\leq C_4 \left(|u - v|^2 + (\varepsilon \theta)^2 |t - s|^2 \right)^{(1-d)/2} \\ &\leq C_5 \left(|u - v| + \varepsilon \theta |t - s| \right)^{1-d} \\ &\leq C_5 \left(|u - v|^{d-1} + (\varepsilon \theta |t - s|)^{d-1} \right)^{-1}. \end{aligned}$$

Thus, also assertion (a) of the claim is true.

Using the result of the above claim, we can continue with the estimate (A.13). Employing a transition to spherical coordinates and that $\varphi_i^{-1}(\Sigma_1) \subset B(u, 2r_0/c)$, we get for all sufficiently small $\varepsilon > 0$ that

$$\begin{split} \int_{-1}^{1} \int_{\varphi_{i}^{-1}(\Sigma_{1})} \left| x(\theta) - \varphi_{i}(v) - \varepsilon \theta s \nu(\varphi_{i}(v)) \right|^{1-d} \mathrm{d}v \mathrm{d}s \\ & \leq C_{5} \int_{-1}^{1} \int_{B(u,2r_{0}/c)} \left(|u - v|^{d-1} + (\varepsilon \theta |t - s|)^{d-1} \right)^{-1} \mathrm{d}v \mathrm{d}s \\ & = C_{6} \int_{-1}^{1} \int_{0}^{2r_{0}/c} r^{d-2} \left(r^{d-1} + (\varepsilon \theta |t - s|)^{d-1} \right)^{-1} \mathrm{d}r \mathrm{d}s. \end{split}$$

We employ now the substitutions $z := r^{d-1}$ and $\zeta := s - t$. Note that $s, t \in [-1, 1]$ implies $\zeta \in [-2, 2]$. Hence, we get

$$\int_{-1}^{1} \int_{\varphi_{i}^{-1}(\Sigma_{1})} |x(\theta) - \varphi_{i}(v) - \varepsilon \theta s \nu(\varphi_{i}(v))|^{1-d} dv ds$$

$$\leq C_{6} \int_{-1}^{1} \int_{0}^{2r_{0}/c} r^{d-2} \left(r^{d-1} + (\varepsilon \theta |t-s|)^{d-1} \right)^{-1} dr ds$$

$$\leq C_{7} \int_{-2}^{2} \int_{0}^{(2r_{0}/c)^{d-1}} \left(z + (\varepsilon \theta |\zeta|)^{d-1} \right)^{-1} dz d\zeta$$

$$= C_{7} \int_{-2}^{2} \left(\ln \left(\left(\frac{2r_{0}}{c} \right)^{d-1} + (\varepsilon \theta |\zeta|)^{d-1} \right) - \ln \left((\varepsilon \theta |\zeta|)^{d-1} \right) \right) d\zeta.$$

For a fixed $r_0 > 0$ it holds that $\ln\left(\left(\frac{2r_0}{c}\right)^{d-1} + (\varepsilon\theta|\zeta|)^{d-1}\right) \leq C_8$ independent of $\theta \in (0,1)$. Therefore, we obtain

$$\int_{-1}^{1} \int_{\varphi_{i}^{-1}(\Sigma_{1})} \left| x(\theta) - \varphi_{i}(v) - \varepsilon \theta s \nu(\varphi_{i}(v)) \right|^{1-d} \mathrm{d}v \mathrm{d}s \leq C_{9} \int_{-2}^{2} \left(1 - \ln(\varepsilon \theta |\zeta|) \right) \mathrm{d}\zeta$$
$$\leq C_{10} \left(1 + |\ln \varepsilon| + |\ln \theta| \right).$$

Summing up all $i \in I$, we get

$$\mathcal{I}_{1}(x(\theta), \theta, \varepsilon) \leq C_{3} \sum_{i \in I} \int_{-1}^{1} \int_{\varphi_{i}^{-1}(\Sigma_{1})} |x(\theta) - \varphi_{i}(v) - \varepsilon \theta s \nu(\varphi_{i}(v))|^{1-d} \mathrm{d}v \mathrm{d}s$$
$$\leq C_{11} (1 + |\ln \varepsilon| + |\ln \theta|).$$

Thus, using (A.12) and the last estimate, it follows

$$\int_{\Sigma} \int_{-1}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| ds d\sigma(y_{\Sigma}) = \mathcal{I}_{1}(x(\theta), \theta, \varepsilon) + \mathcal{I}_{2}(x(\theta), \theta, \varepsilon) \\ \leq C_{12} \left(1 + |\ln \varepsilon| + |\ln \theta| \right).$$

This is the desired estimate, if $dist(x, \Sigma) \leq r_0$. Integrating the last estimate with respect to θ we finally obtain

$$\begin{split} \int_{0}^{1} \int_{\Sigma} \int_{-1}^{1} \left| \nabla G_{\lambda}(x(\theta) - y_{\Sigma} - \varepsilon \theta s \nu(y_{\Sigma})) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \mathrm{d}\theta \\ & \leq C_{13} \int_{0}^{1} \left(1 + |\ln \varepsilon| + |\ln \theta| \right) \mathrm{d}\theta \\ & \leq C_{14} \left(1 + |\ln \varepsilon| \right), \end{split}$$

which leads to the statement of this lemma.

In the following proposition we state the last two estimates that are needed to prove our main result. They follow from Lemma A.8, that is once applied for the case that $x(\theta)$ is a constant function and once for $x(\theta)$ being linear affine in θ .

Proposition A.9 Let Σ be a C^2 -smooth hypersurface satisfying Hypothesis 2.3 and let $\lambda \in (-\infty, 0)$.

(i) Let $\psi \in L^{\infty}(\Sigma \times (-1, 1))$. Then there exists a constant $C = C(d, \lambda, \Sigma, \psi) > 0$ such that the bound

$$\sup_{x \in \mathbb{R}^d} \int_{\Sigma} \int_{-1}^1 \left| \left(G_\lambda(x - y_\Sigma - \varepsilon s\nu(y_\Sigma)) - G_\lambda(x - y_\Sigma) \right) \psi(y_\Sigma, s) \right| \mathrm{d}s \mathrm{d}\sigma(y_\Sigma) \le C\varepsilon \left(1 + |\ln\varepsilon| \right)$$

holds for all sufficiently small $\varepsilon > 0$.

(ii) Let $\omega, \psi \in L^{\infty}(\Sigma \times (-1, 1))$. Then there exists a constant $C = C(d, \lambda, \Sigma, \omega, \psi) > 0$ such that the bound

$$\sup_{(x_{\Sigma},t)\in\Sigma\times(-1,1)}\int_{\Sigma}\int_{-1}^{1}|\omega(x_{\Sigma},t)(G_{\lambda}(x_{\Sigma}+\varepsilon t\nu(x_{\Sigma})-y_{\Sigma}-\varepsilon s\nu(y_{\Sigma})))-G_{\lambda}(x_{\Sigma}-y_{\Sigma}))\psi(y_{\Sigma},s)|dsd\sigma(y_{\Sigma})\leq C\varepsilon(1+|\ln\varepsilon|)$$

holds for all sufficiently small $\varepsilon > 0$.

Proof. (i) Let $x \in \mathbb{R}^d$ be fixed. We employ (A.10) (with $\tau(x, y_{\Sigma}, s, t) = -s\nu(y_{\Sigma})$ implying $\|\tau\|_{L^{\infty}} = 1$) and find that

$$\begin{split} \int_{\Sigma} \int_{-1}^{1} \left| \left(G_{\lambda}(x - y_{\Sigma} - \varepsilon s\nu(y_{\Sigma})) - G_{\lambda}(x - y_{\Sigma}) \right) \psi(y_{\Sigma}, s) \right| \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ & \leq \varepsilon \|\psi\|_{L^{\infty}} \int_{\Sigma} \int_{-1}^{1} \int_{0}^{1} \left| \nabla G_{\lambda}(x - y_{\Sigma} - \varepsilon \theta s\nu(y_{\Sigma})) \right| \mathrm{d}\theta \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}). \end{split}$$

Thus, Lemma A.8 yields the claimed assertion.

(ii) Let $(x_{\Sigma}, t) \in \Sigma \times (-1, 1)$ be fixed. Using (A.10) (with $x = x_{\Sigma}$ and $\tau(x_{\Sigma}, y_{\Sigma}, s, t) = t\nu(x_{\Sigma}) - s\nu(y_{\Sigma})$ implying $\|\tau\|_{L^{\infty}} \leq 2$) we find that

$$\begin{split} \int_{\Sigma} \int_{-1}^{1} |\omega(x_{\Sigma}, t) \big(G_{\lambda}(x_{\Sigma} + \varepsilon t \nu(x_{\Sigma}) - y_{\Sigma} - \varepsilon s \nu(y_{\Sigma})) - G_{\lambda}(x_{\Sigma} - y_{\Sigma}) \big) \psi(y_{\Sigma}, s) | \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}) \\ & \leq 2\varepsilon \|\omega\|_{L^{\infty}} \|\psi\|_{L^{\infty}} \int_{\Sigma} \int_{-1}^{1} \int_{0}^{1} |\nabla G_{\lambda}(x_{\Sigma} - y_{\Sigma} + \varepsilon \theta(t \nu(x_{\Sigma}) - s \nu(y_{\Sigma}))) | \mathrm{d}\theta \mathrm{d}s \mathrm{d}\sigma(y_{\Sigma}). \end{split}$$

Therefore, the claimed statement follows from Lemma A.8 (with $x(\theta) = x_{\Sigma} + \varepsilon \theta t \nu(x_{\Sigma})$).

B Boundaries of bounded C²-domains

In this appendix it is shown that the boundary of a bounded and simply connected C^2 -domain is a C^2 -hypersurface in the sense of Definition 2.1 that satisfies Hypothesis 2.3. Recall that a bounded C^2 -domain $\Omega \subset \mathbb{R}^d$ is an open, bounded and simply connected set with boundary Σ for which there exists a parametrization $\{\Phi_i, X_i, Y_i\}_{i \in I}$ with a finite index set I satisfying the following conditions:

- (a) $X_i \subset \mathbb{R}^{d-1}$ and $Y_i \subset \mathbb{R}^d$ are open sets and $\Phi_i \colon X_i \to Y_i$ is a C^2 -mapping for all $i \in I$;
- (b) rank $D\Phi_i(u) = d 1$ for all $u \in X_i$ and $i \in I$;
- (c) $\Phi_i(X_i) = Y_i \cap \Sigma$ and $\Phi_i \colon X_i \to Y_i \cap \Sigma$ is a homeomorphism;
- (d) $\Sigma \subset \bigcup_{i \in I} Y_i;$

see, e.g. in [30, Chapter 3] for a similar notion. Since Σ is compact it is no restriction to assume that the sets $Y_i, i \in I$, are bounded. In order to prove that Σ is a hypersurface in the sense of Definition 2.1 choose a family $\{V_i\}_{i \in I}$ of open sets in \mathbb{R}^d such that $V_i \subset \overline{V_i} \subset Y_i$ and $\Sigma \subset \bigcup_{i \in I} V_i$. Next, set $U_i := \Phi_i^{-1}(V_i \cap \Sigma) \subset X_i$ and $\varphi_i := \Phi_i|_{U_i}$. Then $U_i \subset \overline{U_i} \subset X_i$ and $\{\varphi_i, U_i, V_i\}_{i \in I}$ satisfies the conditions (a)-(d) above. In addition, each φ_i is uniformly Lipschitz continuous on U_i , as $\overline{U_i} \subset X_i$. Therefore, $\{\varphi_i, U_i, V_i\}_{i \in I}$ is a parametrization of Σ with the properties (a)-(e) in Definition 2.1 and hence Σ is a C^2 -hypersurface.

Moreover, as $\overline{U_i} \subset X_i$ there is a constant C > 0 such that the corresponding first fundamental form in (2.1) satisfies

$$\det G_i(u) \ge C > 0 \tag{B.1}$$

for any $u \in U_i$ and all $i \in I$. Since the outward unit normal vector field ν is continuous, it is clear that the hypersurface Σ is orientable.

Recall that the eigenvalues of the Weingarten map depend continuously on $x_{\Sigma} \in \Sigma$ and, since Σ is compact, these eigenvalues are bounded. Hence, item (b) of Hypothesis 2.3 is satisfied. It remains to show that (a) and (c) of Hypothesis 2.3 hold. The following proposition concerns condition (a).

Proposition B.1 Let $\Omega \subset \mathbb{R}^d$ be a bounded and simply connected C^2 -domain with boundary Σ . Then there exists $\beta > 0$ such that the mapping $\iota_{\Sigma} : \Sigma \times \mathbb{R} \to \mathbb{R}^d$ in (2.5) is injective on $\Sigma \times (-\beta, \beta)$.

Proof. Suppose that the claim is false. Then for all $n \in \mathbb{N}$ there exist $x_{\Sigma,n}, y_{\Sigma,n} \in \Sigma$ and $s_n, t_n \in \left(-\frac{1}{n}, \frac{1}{n}\right)$ such that $(x_{\Sigma,n}, t_n) \neq (y_{\Sigma,n}, s_n)$ and

$$x_{\Sigma,n} + t_n \nu(x_{\Sigma,n}) = \iota_{\Sigma}(x_{\Sigma,n}, t_n) = \iota_{\Sigma}(y_{\Sigma,n}, s_n) = y_{\Sigma,n} + s_n \nu(y_{\Sigma,n}).$$
(B.2)

Since Σ is compact we can assume that the sequences $(x_{\Sigma,n})_n$ and $(y_{\Sigma,n})_n$ converge to x_{Σ} and y_{Σ} , respectively. Then, equation (B.2) implies $x_{\Sigma} = y_{\Sigma}$. Let $\{\varphi_i, U_i, V_i\}_{i \in I}$ be the parametrization of Σ constructed above and let $\iota_{\varphi_i} : U_i \times \mathbb{R} \to \mathbb{R}^d$ be as in (2.3). Since $x_{\Sigma,n} \to x_{\Sigma}$ and $y_{\Sigma,n} \to y_{\Sigma} = x_{\Sigma}$, there exists $N \in \mathbb{N}$ and $i \in I$ such that $x_{\Sigma,n}, y_{\Sigma,n} \in V_i$ for all $n \geq N$. Hence, $u_n := \varphi_i^{-1}(x_{\Sigma,n})$ and $v_n := \varphi_i^{-1}(y_{\Sigma,n}), n \geq N$, satisfy

$$\iota_{\varphi_i}(u_n, t_n) = \varphi_i(u_n) + t_n \nu(\varphi_i(u_n)) = \varphi_i(v_n) + s_n \nu(\varphi_i(v_n)) = \iota_{\varphi_i}(v_n, s_n)$$

and it follows that ι_{φ_i} is not injective on $U_i \times (-\beta, \beta)$ for any $\beta > 0$. On the other hand, if G_i and L_i denote the matrices representing the first fundamental form and the Weingarten map, then

$$|\det D\iota_{\varphi_i}(u,t)| = \det(1 - tL_i(u))\sqrt{\det G_i(u)} \ge C > 0$$

holds for some C > 0, all $u \in U_i$ and all t sufficiently small; cf. (2.4) and (B.1). Now the inverse function theorem implies that ι_{φ_i} is invertible; a contradiction.

Finally, we prove that condition (c) of Hypothesis 2.3 is satisfied for the boundary of a compact and simply connected C^2 -smooth domain.

Proposition B.2 Let $\Omega \subset \mathbb{R}^d$ be a bounded and simply connected C^2 -smooth domain with boundary Σ , let $\{\varphi_i, U_i, V_i\}_{i \in I}$ be the parametrization of Σ as above and let ι_{φ_i} be as in (2.3). Then there exist $\beta > 0$ and a constant c > 0 such that for all $u, v \in U_i$, $s, t \in (-\beta, \beta)$ and all $i \in I$ it holds

$$|\iota_{\varphi_i}(u,t) - \iota_{\varphi_i}(v,s)|^2 \ge c^2 (|u-v|^2 + |s-t|^2).$$

Proof. Let $\{\Phi_i, X_i, Y_i\}_{i \in I}$ be a parametrization of Σ as in the beginning of this appendix, fix $i \in I$ and choose $\beta > 0$ and a constant D > 0 such that

$$\det(1 - tL_i(u)) \ge D > 0 \quad \text{for all} \quad (u, t) \in X_i \times (-2\beta, 2\beta); \tag{B.3}$$

here $L_i(u)$ is the matrix representing the Weingarten map associated to the mapping Φ_i . As in (2.3) define

$$\iota_{\Phi_i} \colon X_i \times \mathbb{R} \to \mathbb{R}^d, \qquad \iota_{\Phi_i}(u, t) := \Phi_i(u) + t\nu(\Phi_i(u)),$$

and consider the open sets

$$\Omega_1 := \iota_{\varphi_i}(U_i \times (-\beta, \beta)) \quad \text{and} \quad \Omega_3 := \iota_{\Phi_i}(X_i \times (-2\beta, 2\beta))$$

in \mathbb{R}^d . Since $\Omega_1 \subset \overline{\Omega_1} \subset \Omega_3$, there exists $\beta_1 > 0$ such that $\overline{B(x, 2\beta_1)} \subset \Omega_3$ for any $x \in \Omega_1$. Eventually, we need the open set

$$\Omega_2 := \bigcup_{x \in \Omega_1} B(x, \beta_1).$$

Note that $\Omega_1 \subset \Omega_2 \subset \overline{\Omega_2} \subset \Omega_3$.

Let $u, v \in U_i$ and $s, t \in (-\beta, \beta)$ be fixed. Set $x := \iota_{\varphi_i}(u, t) = \iota_{\Phi_i}(u, t)$ and $y := \iota_{\varphi_i}(v, s) = \iota_{\Phi_i}(v, s)$. We distinguish two cases: $|x - y| < \beta_1$ and $|x - y| \ge \beta_1$. In the first case, if $|x - y| < \beta_1$ then y is contained in the convex set $B(x, \beta_1) \subset \Omega_2$. It follows from

$$\left|\det\left((D\iota_{\Phi_i})(u,t)\right)\right| = \left|\det(1 - tL_i(u))\right| \sqrt{\det G_i(s)}$$

(cf. (2.4)), (B.3) and (B.1), which is also true for $u \in \Omega_2$ with a possibly smaller C, that $|\det D\iota_{\Phi_i}(u,t)| \ge D' > 0$ for all $(u,t) \in \Phi_i^{-1}(\Omega_2)$. Hence, there exists a constant $C_1 > 0$ such that

$$\left\| (D\iota_{\Phi_{i}}^{-1})(z) \right\| = \left\| \left(D\iota_{\Phi_{i}}(\iota_{\Phi_{i}}^{-1}(z)) \right)^{-1} \right\| = \left\| \frac{1}{\det D\iota_{\Phi_{i}}(\iota_{\Phi_{i}}^{-1}(z))} \operatorname{adj} \left[D\iota_{\Phi_{i}}(\iota_{\Phi_{i}}^{-1}(z)) \right] \right\| \le C_{1}$$

for all $z \in \Omega_2$; here adj $[\cdot]$ denotes the adjugate matrix. Hence, we obtain

$$\begin{aligned} \left(|u-v|^{2}+|s-t|^{2}\right)^{1/2} &= \left|\iota_{\Phi_{i}}^{-1}(y)-\iota_{\Phi_{i}}^{-1}(x)\right| \\ &= \left|\int_{0}^{1} \frac{\mathrm{d}}{\mathrm{d}\xi}\iota_{\Phi_{i}}^{-1}(x+\xi(y-x))\mathrm{d}\xi\right| \\ &\leq \int_{0}^{1} \left|D\iota_{\Phi_{i}}^{-1}(x+\xi(y-x))\cdot(y-x)\right|\mathrm{d}\xi \\ &\leq \max\left\{\left\|D\iota_{\Phi_{i}}^{-1}(z)\right\|: z\in B(x,\beta_{1})\right\}\cdot|x-y| \\ &\leq C_{1}|\iota_{\varphi_{i}}(u,t)-\iota_{\varphi_{i}}(v,s)|, \end{aligned}$$

which is the claimed result in the case $|x - y| < \beta_1$.

In the remaining case $|x - y| \ge \beta_1$ we set $C_2 := \max\{|\iota_{\varphi_i}^{-1}(z)| : z \in \Omega_1\}$. Since $\overline{\Omega}_1 \subset \Omega_3$ we conclude $C_2 < \infty$ and

$$|u-v|^2 + |s-t|^2 = \left|\iota_{\varphi_i}^{-1}(x) - \iota_{\varphi_i}^{-1}(y)\right|^2 \le \frac{4C_2^2}{|x-y|^2}|x-y|^2 \le \frac{4C_2^2}{\beta_1^2}\left|\iota_{\varphi_i}(u,t) - \iota_{\varphi_i}(v,s)\right|^2.$$

Setting finally

$$C^{(i)} := \min\left\{\frac{1}{C_1}, \frac{\beta_1}{2C_2}\right\}$$

and $c := \min \{ C^{(i)} : i \in I \}$, the result of this proposition follows.

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References

- M. Abramowitz and I. Stegun, Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables (Wiley, Washington, D.C., 1984).
- [2] S. Albeverio, F. Gesztesy, R. Høegh-Krohn, and H. Holden, Solvable Models in Quantum Mechanics. With an Appendix by Pavel Exner. 2nd ed. (AMS Chelsea Publishing, Providence, RI, 2005).
- [3] S. Albeverio and V. Koshmanenko, On form-sum approximations of singularly perturbed positive self-adjoint operators, J. Funct. Anal. 169, 32–51 (1999).
- [4] S. Albeverio, A. Kostenko, M. Malamud, and H. Neidhardt, Spherical Schrödinger operators with δ -type interactions, J. Math. Phys. **54**, 052103, 24 pp. (2013).
- [5] J.-P. Antoine, F. Gesztesy, and J. Shabani, Exactly solvable models of sphere interactions in quantum mechanics, J. Phys. A 20, 3687–3712 (1987).
- [6] J. Behrndt, P. Exner, and V. Lotoreichik, Schrödinger operators with δ and δ' -interactions on Lipschitz surfaces and chromatic numbers of associated partitions, Rev. Math. Phys. **26**, 1450015, 43 pp. (2014).
- [7] J. Behrndt, P. Exner, and V. Lotoreichik, Schrödinger operators with δ -interactions supported on conical surfaces, J. Phys. A **47**, 355202, 16 pp. (2014).
- [8] F.A. Berezin and L.D. Faddeev, A remark on Schrödinger's equation with a singular potential, Sov. Math., Dokl. 2, 372–375 (1961).
- [9] J. Brasche, On the spectral properties of singular perturbed operators, in: Dirichlet forms and stochastic processes. Proceedings of the international conference held in Beijing, China, October 25-31, 1993, (Berlin: de Gruyter, 1995), pp. 65–72.
- [10] J. Brasche, P. Exner, Y. Kuperin, and P. Šeba, Schrödinger operators with singular interactions, J. Math. Anal. Appl. 184, 112–139 (1994).
- [11] J. Brasche, R. Figari, and A. Teta, Singular Schrödinger operators as limits of point interaction Hamiltonians, Potential Anal. 8, 163–178 (1998).
- [12] J. Brasche and K. Ožanová, Convergence of Schrödinger operators, SIAM J. Math. Anal. 39, 281–297 (2007).
- [13] M. P. do Carmo, Differential Geometry of Curves and Surfaces (Englewood Cliffs, N.J., 1976).
- [14] P. Duclos, P. Exner, and D. Krejčiřík, Bound states in curved quantum layers, Comm. Math. Phys. 223, 13–28 (2001).
- [15] P. Exner, Leaky quantum graphs: a review, in: Analysis on graphs and its applications. Selected papers based on the Isaac Newton Institute for Mathematical Sciences programme, Cambridge, UK, January 8–June 29, 2007, (Providence, RI: American Mathematical Society (AMS), 2008), pp. 523–564.
- [16] P. Exner and T. Ichinose, Geometrically induced spectrum in curved leaky wires, J. Phys. A, Math. Gen. 34, 1439–1450 (2001).
- [17] P. Exner and K. Yoshitomi, Band gap of the Schrödinger operator with a strong δ -interaction on a periodic curve, Ann. Henri Poincaré **2**, 1139–1158 (2001).
- [18] P. Exner and S. Kondej, Bound states due to a strong δ interaction supported by a curved surface, J. Phys. A **36**, 443–457 (2003).
- [19] P. Exner and H. Kovařík, Quantum Waveguides (Cham, Springer, 2015),
- [20] P. Exner and K. Němcová, Leaky quantum graphs: Approximations by point-interaction Hamiltonians, J. Phys. A 36, 10173–10193 (2003).
- [21] A. Figotin and P. Kuchment, Spectral properties of classical waves in high-contrast periodic media, SIAM J. Appl. Math. 58, 683–702 (1998).
- [22] G. E. Galkowski, Distribution of Resonances in Scattering by Thin Barriers (UC Berkeley Electronic Theses and Dissertations, 2015).
- [23] D. Haroske and H. Triebel, Distributions, Sobolev spaces, Elliptic Equations (EMS Textbooks in Mathematics, Zürich, 2008).
- [24] A. Jonsson and H. Wallin, Function Spaces on Subsets of \mathbb{R}^n (Math. Rep. Ser. 2, 1984).
- [25] T. Kato, Perturbation Theory for Linear Operators (Springer-Verlag, Berlin, 1995).
- [26] R. de L. Kronig and W. Penney, Quantum mechanics of electrons in crystal lattices, Proc. Roy. Soc. Lond. 130, 499–513 (1931).
- [27] W. Kühnel, Differential Geometry. Curves-surfaces-manifolds (American Mathematical Society, Providence, RI, 2002).
- [28] G. Leoni, A First Course in Sobolev Spaces (American Mathematical Society, Providence, RI, 2009).
 [29] Ch. Lin and Z. Lu, Existence of bound states for layers built over hypersurfaces in ℝⁿ⁺¹, J. Funct. Anal. 244, 1–25 (2007).
- [30] W. McLean, Strongly Elliptic Systems and Boundary Integral Equations (Cambridge University Press, Cambridge, 2000).
- [31] K. Ozanová, Approximation by point potentials in a magnetic field, J. Phys. A 39, 3071-3083 (2006).
- [32] I. Yu. Popov, The operator extension theory, semitransparent surface and short range potential, Math. Proc. Camb. Phil. Soc. 118, 555–563 (1995).
- [33] N. Setô, Bargmann's inequalities in spaces of arbitrary dimension, Publ. Res. Inst. Math. Sci. 9, 429–461 (1974).

- [34] S. Shimada, The approximation of the Schrödinger operators with penetrable wall potentials in terms of short range Hamiltonians, J. Math. Kyoto Univ. **32**, 583–592 (1992).
- [35] P. Stollmann and J. Voigt, Perturbation of Dirichlet forms by measures, Potential Anal. 5, 109–138 (1996).
- [36] G. Teschl, Mathematical Methods in Quantum Mechanics. With Applications to Schrödinger Operators. (American Mathematical Society, Providence, 2014).
- [37] L. H. Thomas, The interaction between a neutron and a proton and the structure of H^3 , Phys. Rev., II. Ser. 47, 903–909 (1935).
- [38] J. Weidmann, Lineare Operatoren in Hilberträumen. Teil 1. (Teubner, Stuttgart, 2000).
- [39] K. Yoshitomi, Band gap of the spectrum in periodically curved quantum waveguides, J. Differ. Equations **142**, 123–166 (1998).