# Selfadjoint Schrödinger operators on the half-space with compactly supported Robin boundary conditions

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#### 1 Introduction

We investigate realizations of the differential expression  $-\Delta + V$  on the half-space  $\mathbb{R}^n_+ = \{(x', x_n) : x' \in \mathbb{R}^{n-1}, x_n > 0\}, n \ge 2$ , with a real-valued, bounded potential V. More precisely, we study the differential operator

$$A_g u = -\Delta u + V u, \quad \operatorname{dom} A_g = \left\{ u \in H^{3/2}_{\Delta}(\mathbb{R}^n_+) : \partial_{\nu} u|_{\mathbb{R}^{n-1}} = g \cdot (u|_{\mathbb{R}^{n-1}}) \right\},\tag{1}$$

in  $L^2(\mathbb{R}^n_+)$ , where  $H^{3/2}_{\Delta}(\mathbb{R}^n_+) = \{u \in H^{3/2}(\mathbb{R}^n_+) : \Delta u \in L^2(\mathbb{R}^n_+)\}$  and  $g : \mathbb{R}^{n-1} \to \mathbb{R}$  is a bounded, real function with compact support. The aim of the present note is to show that  $A_g$  is a selfadjoint, compact perturbation in the resolvent sense of the selfadjoint realization  $A_N$  of  $-\Delta + V$  with Neumann boundary conditions. In particular this guarantees that  $A_g$  and  $A_N$  have the same essential spectrum. We point out that the latter can still be proved under slightly weaker assumptions on g, see [10] for a more general approach and [8] for a result with a more regular g in dimension n = 2. Our proofs make use of techniques which were originally developed in [2, 3] for the treatment of elliptic differential operators on domains with a compact boundary. For further recent developments in this area we refer the reader to [1,4,6,11,12].

#### 2 Preliminaries

In this section we fix some notation and recall some known facts on Sobolev spaces and Schrödinger operators; proofs and further details can be found in [9] and, e.g., [7, Chapter 9]. Let  $K \subset \mathbb{R}^{n-1}$  be a compact set and let  $H^s(\mathbb{R}^n_+)$  and  $H^s(K) = \{f|_K : f \in H^s(\mathbb{R}^{n-1})\}$  be the Sobolev spaces of order s > 0 on  $\mathbb{R}^n_+$  and K, respectively. For  $u \in H^{3/2}_{\Delta}(\mathbb{R}^n_+)$  we denote by  $u|_{\mathbb{R}^{n-1}}$  the trace of u on the boundary  $\mathbb{R}^{n-1}$  of  $\mathbb{R}^n_+$ , by  $\partial_{\nu} u|_{\mathbb{R}^{n-1}} = -\frac{\partial u}{\partial x_n}|_{\mathbb{R}^{n-1}}$  the derivative of u along the outer normal vector field on  $\mathbb{R}^{n-1}$ , and by  $u|_K$  and  $\partial_{\nu} u|_K$ ,  $\partial_{\nu} u|_{\mathbb{R}^{n-1}\setminus K}$  their restrictions to K and  $\mathbb{R}^{n-1} \setminus K$ , respectively. The mappings  $\Gamma_0$  and  $\Gamma_1$  given by

$$\Gamma_0: H^{3/2}_{\Delta}(\mathbb{R}^n_+) \to L^2(K), \quad \Gamma_0 u = \partial_{\nu} u|_K \quad \text{and} \quad \Gamma_1: H^{3/2}_{\Delta}(\mathbb{R}^n_+) \to H^1(K), \quad \Gamma_1 u = u|_K$$
(2)

are surjective.

Here and in the following let  $V \in L^{\infty}(\mathbb{R}^n_+)$  be real-valued. It is well known that the *Neumann operator* 

$$A_N u = -\Delta u + V u, \qquad \operatorname{dom} A_N = \left\{ u \in H^{3/2}_{\Delta}(\mathbb{R}^n_+) : \partial_{\nu} u|_{\mathbb{R}^{n-1}} = 0 \right\}$$

is a selfadjoint realization of  $-\Delta + V$  in  $L^2(\mathbb{R}^n_+)$ , and by elliptic regularity dom  $A_N \subset H^2(\mathbb{R}^n_+)$  holds. Note that this yields the decomposition  $\{u \in H^{3/2}_{\Delta}(\mathbb{R}^n_+) : \partial_{\nu}u|_{\mathbb{R}^{n-1}\setminus K} = 0\} = \operatorname{dom} A_N + \mathcal{N}_{\lambda}$  for each  $\lambda$  in the resolvent set  $\rho(A_N)$  of  $A_N$ , where  $\mathcal{N}_{\lambda} := \{u \in H^{3/2}(\mathbb{R}^n_+) : -\Delta u + Vu = \lambda u, \partial_{\nu}u|_{\mathbb{R}^{n-1}\setminus K} = 0\}$ . This, together with (2), ensures that the *Poisson* operator

$$\gamma(\lambda): L^2(K) \to L^2(\mathbb{R}^n_+), \quad \partial_\nu u_\lambda|_K \mapsto u_\lambda, \quad u_\lambda \in \mathcal{N}_\lambda, \tag{3}$$

and the Neumann-to-Dirichlet operator

$$M(\lambda): L^2(K) \to L^2(K), \quad \partial_\nu u_\lambda|_K \mapsto u_\lambda|_K, \quad u_\lambda \in \mathcal{N}_\lambda, \tag{4}$$

are well-defined for each  $\lambda \in \rho(A_N)$ . Moreover,  $\gamma(\lambda)$  and  $M(\lambda)$  are bounded and ran  $M(\lambda) = H^1(K)$  holds.

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## 3 Selfadjoint Schrödinger operators on the half-space

The following theorem is the main result of this note. For  $g \in L^{\infty}(\mathbb{R}^{n-1})$ , supp g = K, we denote by G the operator of multiplication with the function  $g|_K$  in  $L^2(K)$ .

**Theorem 3.1** Let  $K \subset \mathbb{R}^{n-1}$  be a compact set and let  $g \in L^{\infty}(\mathbb{R}^{n-1})$  be a real-valued function with supp g = K. Then the operator  $A_g$  in (1) is selfadjoint in  $L^2(\mathbb{R}^n_+)$  and  $\lambda \in \rho(A_N)$  is an eigenvalue of  $A_g$  if and only if 1 is an eigenvalue of  $GM(\lambda)$ . The resolvent difference

$$(A_g - \lambda)^{-1} - (A_N - \lambda)^{-1} = \gamma(\lambda) \left( I - GM(\lambda) \right)^{-1} G\gamma(\overline{\lambda})^*, \quad \lambda \in \rho(A_g) \cap \rho(A_N),$$
(5)

is compact and, in particular, the essential spectra of  $A_g$  and  $A_N$  coincide.

Proof. Let us first show that  $\lambda \in \rho(A_N)$  is an eigenvalue of  $A_g$  if and only if 1 is an eigenvalue of  $GM(\lambda)$ . For  $u \in \ker(A_g - \lambda), u \neq 0$ , we have  $\Gamma_0 u \neq 0$  and  $GM(\lambda)\Gamma_0 u = G\Gamma_1 u = \Gamma_0 u$ . Thus  $I - GM(\lambda)$  is not injective. Conversely,  $f \in \ker(I - GM(\lambda)), f \neq 0$ , implies  $\gamma(\lambda)f \in \operatorname{dom} A_g, (A_g - \lambda)\gamma(\lambda)f = 0$ , and  $\gamma(\lambda)f \neq 0$ . Thus  $\gamma(\lambda)f$  is an eigenfunction of  $A_g$  corresponding to the eigenvalue  $\lambda$ .

Next we show that  $A_g$  is a selfadjoint operator in  $L^2(\mathbb{R}^n_+)$ . For this note first that for  $u \in \text{dom } A_q$  we have

$$(A_g u, u) = \int_{\mathbb{R}^n_+} (-\Delta + V) u \,\overline{u} \, dx = \int_{\mathbb{R}^n_+} |\nabla u|^2 + V |u|^2 \, dx - \int_K g \, |u|^2 \, d\sigma \, \in \, \mathbb{R},$$

so that  $A_g$  is a symmetric in operator in  $L^2(\mathbb{R}^n_+)$ . Hence it is sufficient to verify that  $A_g - \lambda$  is surjective for  $\lambda \in \mathbb{C} \setminus \mathbb{R}$ . Fix some  $\lambda \in \mathbb{C} \setminus \mathbb{R}$ , choose an arbitrary  $u \in L^2(\mathbb{R}^n_+)$ , and define

$$v := (A_N - \lambda)^{-1} u + \gamma(\lambda) \left( I - GM(\lambda) \right)^{-1} G\gamma(\overline{\lambda})^* u.$$
(6)

In the following we will show that v is well-defined and belongs to dom  $A_g$  with  $(A_g - \lambda)v = u$ . The operator  $\gamma(\lambda)$  and hence also  $\gamma(\overline{\lambda})^*$  and  $G\gamma(\overline{\lambda})^*$  are bounded and everywhere defined. Furthermore, since  $\operatorname{ran} M(\lambda) = H^1(K)$  and the embedding from  $H^1(K)$  into  $L^2(K)$  is compact,  $M(\lambda)$  and  $GM(\lambda)$  are also compact operators in  $L^2(K)$ . Together with the fact that 1 is not an eigenvalue of  $GM(\lambda)$  we conclude that the operator  $I - GM(\lambda)$  has an everywhere defined, bounded inverse, i.e., v in (6) is well-defined. From the definition of v it is easy to see that  $v \in H^{3/2}_{\Delta}(\mathbb{R}^n_+)$  and  $\partial_{\nu}v|_{\mathbb{R}^{n-1}\setminus K} = 0$  holds.

It remains to show  $G\Gamma_1 v = \Gamma_0 v$  and  $(A_g - \lambda)v = u$ . In fact, as a consequence of the second Green identity we find  $\Gamma_1(A_N - \lambda)^{-1}u = \gamma(\overline{\lambda})^*u$  and therefore we conclude from (6)

$$G\Gamma_1 v = G\gamma(\overline{\lambda})^* u + GM(\lambda) \left( I - GM(\lambda) \right)^{-1} G\gamma(\overline{\lambda})^* u = \left( I - GM(\lambda) \right)^{-1} G\gamma(\overline{\lambda})^* u = \Gamma_0 v.$$

Thus we have shown  $v \in \text{dom } A_g$  and from  $(A_g - \lambda)v = (-\Delta + V - \lambda)v = u$  we obtain that  $A_g - \lambda$  is surjective and, hence,  $A_g$  is selfadjoint. Moreover, we have shown the formula (5) for all  $\lambda \in \mathbb{C} \setminus \mathbb{R}$  and the same reasoning applies for real  $\lambda \in \rho(A_N) \cap \rho(A_g)$ . As mentioned above,  $\gamma(\overline{\lambda})^* = \Gamma_1(A_N - \lambda)^{-1}$ , in particular,  $\operatorname{ran} \gamma(\overline{\lambda})^* \subset H^{3/2}(K)$ , which is compactly embedded in  $L^2(K)$ . This shows that the right hand side in (5) is compact for  $\lambda \in \mathbb{C} \setminus \mathbb{R}$ . Hence  $(A_g - \lambda)^{-1} - (A_N - \lambda)^{-1}$ is compact for each  $\lambda \in \rho(A_g) \cap \rho(A_N)$ , and, in particular,  $A_g$  and  $A_N$  have the same essential spectrum.

We obtain the following corollary in the case V = 0.

**Corollary 3.2** Let V = 0. Then the essential spectrum of the operator  $A_g$  in (1) is given by  $[0, +\infty)$ . Moreover,  $\lambda < 0$  is an eigenvalue of  $A_g$  if and only if 1 is an eigenvalue of  $G\iota^*(-\Delta_{\mathbb{R}^{n-1}} - \lambda)^{-1/2}\iota$ , where  $\iota$  denotes the embedding from  $L^2(K)$  into  $L^2(\mathbb{R}^{n-1})$  and  $\Delta_{\mathbb{R}^{n-1}}$  is the Laplacian on  $\mathbb{R}^{n-1}$ .

Proof. In the case V = 0 it is well-known that the spectrum and essential spectrum of  $A_N$  is given by  $[0, +\infty)$ . Moreover, one computes similarly as in [7, Chapter 9] that  $M(\lambda) = \iota^*(-\Delta_{\mathbb{R}^{n-1}} - \lambda)^{-1/2}\iota$  holds.

### References

- [1] H. Abels, G. Grubb, and I. G. Wood, submitted.
- [2] J. Behrndt and M. Langer, J. Funct. Anal. 243 (2007), 536–565.
- [3] J. Behrndt, M. Langer, and V. Lotoreichik, submitted.
- [4] B. M. Brown, G. Grubb, and I. G. Wood, Math. Nachr. 282 (2009), 314–347.
- [5] F. Gesztesy and M. Mitrea, J. Differential Equations 247 (2009), 2871–2896.
- [6] F. Gesztesy and M. Mitrea, J. Anal. Math. **113** (2011), 53–172.
- [7] G. Grubb, Distributions and Operators (Springer, New York, 2009).
- [8] M. Jílek, SIGMA Symmetry Integrability Geom. Methods Appl. 3 (2007), 12 pp.
- [9] J. L. Lions and E. Magenes, Non-Homogeneous Boundary Value Problems and Applications I (Springer, Berlin, 1972).
- [10] V. Lotoreichik and J. Rohleder, submitted.
- [11] M. M. Malamud, Russ. J. Math. Phys. 17 (2010), 96-125.
- [12] A. Posilicano and L. Raimondi, J. Phys. A: Math. Theor. 42 (2009), 11 pp.