# Eigenvalues of Schrödinger operators and Dirichlet-to-Neumann maps

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Eigenvalues and eigenspaces of selfadjoint Schrödinger operators on  $\mathbb{R}^n$  are expressed in terms of Dirichlet-to-Neumann maps corresponding to Schrödinger operators on the upper and lower half space.

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## 1 Introduction

It is known that the eigenvalues of a Schrödinger operator  $A_D$  with Dirichlet boundary condition on a bounded domain  $\Omega \subset \mathbb{R}^n$  with a bounded, real-valued potential V coincide with the poles of the meromorphic operator function  $\mu \mapsto M^{\Omega}(\mu)$ , where  $M^{\Omega}(\mu)$  is the Dirichlet-to-Neumann map of  $-\Delta + V - \mu$ , see, e.g., [1,2]. Moreover, for each eigenvalue  $\lambda$  the map

 $\tau: \ker(A_D - \lambda) \to \operatorname{ran} \operatorname{Res}_{\lambda} M^{\Omega}, \qquad u \mapsto \partial_{\nu} u|_{\partial\Omega}$ 

(where  $\partial_{\nu}u|_{\partial\Omega}$  denotes the trace of the normal derivative of u at the boundary  $\partial\Omega$ ) is an isomorphism between the eigenspace and the range of the residue of  $M^{\Omega}$  at  $\lambda$ ; cf. [2]. Such a result is also desirable for a selfadjoint Schrödinger operator  $A = -\Delta + V$  in  $L^2(\mathbb{R}^n)$ ,  $n \geq 2$ . In order to define an operator function which plays the role of  $M^{\Omega}$  we introduce the artificial "boundary"  $\Sigma := \mathbb{R}^{n-1} \times \{0\}$ , which separates  $\mathbb{R}^n$  into  $\mathbb{R}^n_+ := \mathbb{R}^{n-1} \times (0, \infty)$  and  $\mathbb{R}^n_- := \mathbb{R}^{n-1} \times (-\infty, 0)$ , and consider the Dirichlet-to-Neumann maps  $M^{\pm}(\mu)$  in  $L^2(\Sigma)$  corresponding to the Schrödinger operators  $-\Delta + V - \mu$  on  $\mathbb{R}^n_{\pm}$ , respectively. A natural candidate for the description of the eigenvalues of A is  $M(\mu) := (M^+(\mu) + M^-(\mu))^{-1}$ ; cf. [3] for a similar function defined in the case that  $\Sigma$  is a sphere. In Theorem 2.1 of this note we show that each pole of M is an eigenvalue of A but in general the analog of the map  $\tau$  is not bijective. We indicate in Theorem 2.2 that this drawback can be avoided by considering a certain  $2 \times 2$  block operator matrix function with entries formed by  $M^{\pm}$  and M.

### 2 Characterization of eigenvalues and eigenspaces with Dirichlet-to-Neumann maps

Let  $n \ge 2$  and denote by  $H^s(\mathbb{R}^n)$  and  $H^s(\Sigma)$  the Sobolev spaces of order s > 0 on  $\mathbb{R}^n$  and  $\Sigma$ , respectively. Moreover, let  $V \in L^{\infty}(\mathbb{R}^n)$  be a real-valued potential. We consider the selfadjoint Schrödinger operator

$$Au = -\Delta u + Vu, \qquad \text{dom} A = H^2(\mathbb{R}^n),$$

in  $L^2(\mathbb{R}^n)$ . For  $\mu$  in the resolvent set  $\rho(A)$  of A we define

$$\mathcal{N}_{\mu}^{\pm} := \{ u_{\mu}^{\pm} \in H^2(\mathbb{R}_{\pm}^n) : (-\Delta + V - \mu) u_{\mu}^{\pm} = 0 \},\\ \mathcal{N}_{\mu} := \{ u_{\mu}^{+} \oplus u_{\mu}^{-} \in \mathcal{N}_{\mu}^{+} \oplus \mathcal{N}_{\mu}^{-} : u_{\mu}^{+} |_{\Sigma} = u_{\mu}^{-} |_{\Sigma} \},$$

where  $v|_{\Sigma}$  denotes the trace of a Sobolev function v at  $\Sigma$ . Let  $\partial_n v := \frac{\partial v}{\partial x_n}$ . One can show, that for every  $g \in H^{\frac{1}{2}}(\Sigma)$  there exists a unique element  $u_{\mu} \in \mathcal{N}_{\mu}$  with  $\partial_n u_{\mu}^-|_{\Sigma} - \partial_n u_{\mu}^+|_{\Sigma} = g$ . Hence the operator-valued function M defined via

$$\rho(A) \ni \mu \mapsto M(\mu), \qquad M(\mu) \left( \partial_n u_{\mu}^{-} |_{\Sigma} - \partial_n u_{\mu}^{+} |_{\Sigma} \right) := u_{\mu}|_{\Sigma}$$

is well-defined.  $M(\mu)$  is a bounded operator in  $L^2(\Sigma)$  with domain  $H^{\frac{1}{2}}(\Sigma)$  and range in  $H^{\frac{3}{2}}(\Sigma)$  for every  $\mu \in \rho(A)$ . Moreover, for every  $g \in H^{\frac{1}{2}}(\Sigma)$  the function  $\mu \mapsto M(\mu)g$  is holomorphic and has poles of at most order one; cf. [2]. Note that for  $\mu \in \mathbb{C} \setminus \mathbb{R}$  the operator  $M(\mu)$  coincides with  $(M^+(\mu) + M^-(\mu))^{-1}$ , where  $M^{\pm}(\mu)$  denotes the Dirichlet-to-Neumann map with respect to  $-\Delta + V - \mu$  on  $\mathbb{R}^n_{\pm}$ , i.e.  $M^{\pm}(\mu)u^{\pm}_{\mu}|_{\Sigma} = \mp \partial_n u^{\pm}_{\mu}|_{\Sigma}$  for  $u^{\pm}_{\mu} \in \mathcal{N}^{\pm}_{\mu}$ , respectively.

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**Theorem 2.1** If  $\lambda \in \mathbb{R}$  is a pole of M then  $\lambda$  is an eigenvalue of A, but in general dim ran  $\operatorname{Res}_{\lambda} M \leq \dim \ker(A - \lambda)$ .

Proof. Let  $\lambda \in \mathbb{R}$  be a pole of M. We show dim ker $(A - \lambda) \geq \dim \operatorname{ran} \operatorname{Res}_{\lambda} M$ , from which, in particular, the first assertion follows. Let  $\mu, \nu, z \in \mathbb{C} \setminus \mathbb{R}$  be distinct and let  $g \in H^{\frac{1}{2}}(\Sigma)$ . For  $j, k \in \{\mu, \nu, z\}$  denote by  $u_j$  the unique element in  $\mathcal{N}_j$  with  $\partial_n u_j^-|_{\Sigma} - \partial_n u_j^+|_{\Sigma} = g$  and choose  $u_k$  analogously. Due to  $u_j - u_k \in \operatorname{dom} A$  and

$$(A-j)(u_j - u_k) = (-\Delta + V - j)(u_j^+ - u_k^+) \oplus (-\Delta + V - j)(u_j^- - u_k^-) = (j-k)u_k$$

we obtain  $(A - j)^{-1}u_k = \frac{u_j - u_k}{j - k}$  if  $j \neq k$ . Hence we get

$$\left( (A-\mu)^{-1} (A-z)^{-1} u_{\nu} \right) \Big|_{\Sigma} = \frac{1}{z-\nu} \left( (A-\mu)^{-1} (u_{z}-u_{\nu}) \right) \Big|_{\Sigma} = \frac{1}{z-\nu} \left[ \frac{u_{\mu}-u_{z}}{\mu-z} - \frac{u_{\mu}-u_{\nu}}{\mu-\nu} \right] \Big|_{\Sigma}$$
$$= \frac{1}{z-\nu} \left[ \frac{M(\mu)g - M(z)g}{\mu-z} - \frac{M(\mu)g - M(\nu)g}{\mu-\nu} \right].$$

By the spectral theorem one gets  $iPu_{\nu} = \lim_{\eta \searrow 0} \eta (A - (\lambda + i\eta))^{-1} u_{\nu}$ , where P denotes the orthogonal projection in  $L^2(\mathbb{R}^n)$  onto  $\ker(A - \lambda)$ . As the map  $v \mapsto [(A - \mu)^{-1}v]|_{\Sigma}$  is continuous from  $L^2(\mathbb{R}^n)$  to  $L^2(\Sigma)$  we get for  $z = \lambda + i\eta$ 

$$\begin{aligned} (Pu_{\nu})|_{\Sigma} &= \left[ (A-\mu)^{-1} (\lambda-\mu) Pu_{\nu} \right]|_{\Sigma} = (\lambda-\mu) \lim_{\eta \searrow 0} \frac{\eta}{i} \left[ (A-\mu)^{-1} (A-(\lambda+i\eta))^{-1} u_{\nu} \right]|_{\Sigma} \\ &= \lim_{\eta \searrow 0} \frac{(\lambda-\mu)\eta}{(z-\nu)i} \left[ \frac{M(\mu)g - M(z)g}{\mu-z} - \frac{M(\mu)g - M(\nu)g}{\mu-\nu} \right] = \lim_{\eta \searrow 0} \frac{i\eta}{\lambda-\nu} M(z)g = \frac{\operatorname{Res}_{\lambda} Mg}{\lambda-\nu}. \end{aligned}$$

We have shown  $\{u|_{\Sigma} : u \in P\mathcal{N}_{\nu}\} = \operatorname{ran} \operatorname{Res}_{\lambda} M$ , hence  $\dim \ker(A - \lambda) \ge \dim \operatorname{ran} \operatorname{Res}_{\lambda} M$ . In general equality does not hold. For example for a potential V reflection symmetric with respect to  $\Sigma$  (i.e.,  $V(x', x_n) = V(x', -x_n)$ ) eigenfunctions with vanishing traces on  $\Sigma$  may exist.

In order to characterize all eigenvalues and eigenspaces of A we define the block operator matrix function  $\mathcal{M}$  via

$$\mu \mapsto \mathcal{M}(\mu) := \begin{bmatrix} M(\mu) & -M(\mu)M^{-}(\mu) \\ -M^{-}(\mu)M(\mu) & -M^{-}(\mu)M(\mu)M^{+}(\mu) \end{bmatrix}, \qquad \mu \in \mathbb{C} \setminus \mathbb{R}.$$

 $\mathcal{M}(\mu)$  is an operator in  $L^2(\Sigma) \times L^2(\Sigma)$  with domain  $H^{\frac{1}{2}}(\Sigma) \times H^{\frac{3}{2}}(\Sigma)$  and range in  $H^{\frac{3}{2}}(\Sigma) \times H^{\frac{1}{2}}(\Sigma)$ . The function  $\mathcal{M}$  is holomorphic in the strong sense and can be extended to a strongly holomorphic function (also denoted by  $\mathcal{M}$ ) defined on  $\rho(A)$ . Similar functions were already considered in, e.g., [5] for the ODE case and in [6,7] in an abstract setting.

**Theorem 2.2**  $\lambda \in \mathbb{R}$  is a pole of  $\mathcal{M}$  and ran  $\operatorname{Res}_{\lambda} \mathcal{M}$  is finite-dimensional if and only if  $\lambda$  is an isolated eigenvalue of A with finite multiplicity. In this case the map

$$\mathcal{T}: \ker(A - \lambda) \to \operatorname{ran} \operatorname{Res}_{\lambda} \mathcal{M}, \qquad u \mapsto [u|_{\Sigma}, -\partial_n u|_{\Sigma}]^{\top}.$$

is bijective.

We omit the proof of Theorem 2.2, which uses methods similar to the proof of Theorem 2.1 and a unique continuation argument; cf. [4] for a similar reasoning.

**Remark 2.3** With the help of the function  $\mathcal{M}$  one can even characterize all (embedded and isolated) eigenvalues and the corresponding eigenspaces of A; cf. [4] for the case of a Schrödinger operator on an exterior domain.

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